

Wave effects on coastal upwelling and water level

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Abstract

Traditional atmosphere, ocean and wave models are run independently of each other. This means that the energy and momentum fluxes do not fully account for the impact of the oceanic wave field at the air-sea interface. In this study, the Stokes drift impact on mass and tracer advection, the Stokes-Coriolis forcing and, the sea-state-dependent momentum and energy fluxes are introduced into an ocean circulation model and tested for a domain covering the Baltic Sea and the North Sea. Sensitivity experiments are designed to investigate the influence on the simulation of storms and Baltic Sea upwelling. Inclusion of wave effects improves the model performance compared with the stand-alone circulation model in terms of sea level height, temperature and circulation. The direct sea-state-dependent momentum and turbulent kinetic energy fluxes prove to be of higher importance than the Stokes drift related effects investigated in this study (i.e., Stokes-Coriolis forcing and Stokes drift advection on tracers and on mass). The latter affects the mass and tracer advection but largely balances the influence of the Stokes-Coriolis forcing. The upwelling frequency changes by more than 10% along the Swedish coast when wave effects are included. In general, the strong (weak) upwelling probability is reduced (increased) when adding the wave effects. From the results, we

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conclude that inclusion of wave effects can be important for regional, high-resolution ocean models even on short time scales, suggesting that they should be introduced in operational ocean circulation models. However, care should be taken when introducing the Stokes-Coriolis forcing as it should be balanced by the Stokes drift in mass and tracer advection.

Keywords: Surface waves; Stokes-Coriolis forcing; Stokes advection; Coupling; Sea-state-dependent fluxes

1. Introduction

Waves play a buffering role in air-sea interaction processes and affect the air-sea interface directly on scales up to the depth of the mixed layer (i.e., Axell, 2002; Sullivan et al., 2007; Sullivan and McWilliams, 2010; McWilliams et al., 2012; Li et al., 2017) and indirectly even deeper (e.g., Breivik et al., 2015). However, the impact of surface waves on air-sea interaction processes is often ignored since the temporal and spatial scales of surface waves are much smaller than the atmospheric and oceanic dynamic scales (Hasselmann, 1991).

The flow of the oceanic and atmospheric boundary layers is affected by surface waves, as has been observed in the open ocean (e.g. Höglström et al., 2015) and laboratory experiments (e.g. Buckley and Veron, 2016) and has been simulated with numerical models (e.g. Sullivan and McWilliams, 2010). Surface waves affect the momentum flux/wind stress (Wu et al., 2016, 2017; Staneva et al., 2017), heat fluxes (Veron et al., 2008) and mixed layer turbulence levels (Craig and Banner, 1994; Sullivan et al., 2004; Sullivan and McWilliams, 2010; Cavaleri et al., 2018). Traditionally, atmospheric and oceanic models have been run separately without a wave model to represent the boundary between the two media. This can cause biases in the upper ocean due to insufficient (or in some cases too strong, see Breivik et al. 2015) mixing and because the momentum transfer is shifted in time and space compared to how the fluxes would behave in the presence of waves. Recently, several sea-state-dependent momentum flux (i.e., Guan and Xie, 2004; Wu et al., 2016) and turbulence closure parameterizations have been proposed (i.e., McWilliams and Sullivan, 2000; Smyth et al., 2002; Qiao

et al., 2004; Huang and Qiao, 2010) and improved performance has been reported for global and regional oceanic and coupled ocean-wave-atmosphere model systems
 25 (Breivik et al., 2015; Staneva et al., 2017; Wu et al., 2015, 2016). Here we will briefly summarize the various wave-induced and wave-related processes that affect the upper ocean before we explain and motivate which processes we have chosen to model.

1.1. Stokes drift related processes

Stokes drift, $\mathbf{u}_S$, is the drift in the wave propagation direction induced by the motion
 30 of surface waves, which can be expressed as the mean temporal and spatial difference between the Eulerian, $\mathbf{u}$, and Lagrangian velocities, $\mathbf{u}_L$ (Stokes, 1847),

$$\mathbf{u}_S = \mathbf{u}_L - \mathbf{u}. \quad (1)$$

Three terms related to the Stokes drift are indicated in the wave-averaged momentum equation, which can be written, loosely following the notation adopted by Suzuki and Fox-Kemper (2016) (see their Eq. 2), as

$$\frac{D\mathbf{u}}{Dt} = -\nabla\left(\frac{p}{\rho_w} + \underbrace{\frac{1}{2}|\mathbf{u}_L|^2 - \frac{1}{2}|\mathbf{u}|^2}_{\text{Press corr}}\right) + \mathbf{u} \times f\hat{\mathbf{z}} + \underbrace{\mathbf{u}_S \times f\hat{\mathbf{z}}}_{\text{CSF}} + \mathbf{D}^u - g\hat{\mathbf{z}} - \underbrace{(\nabla \times \mathbf{u}) \times \mathbf{u}_S}_{\text{Vortex force}}. \quad (2)$$

35 Here p the pressure, $\hat{\mathbf{z}}$ is the upward unit vector, $\mathbf{D}^u$ represents parameterizations of sub-grid scale processes, and $\frac{D\mathbf{u}}{Dt} = \frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla)\mathbf{u}$ is the material or total derivative. The wave effects manifest themselves as the Coriolis-Stokes force (CSF) (Hasselmann, 1970), the vortex force and a Stokes-corrected pressure (McWilliams and Restrepo, 1999).

40 Langmuir turbulence (LT) is believed to influence the whole ocean surface boundary layer (e.g., Sullivan et al., 2007; McWilliams et al., 2012; Belcher et al., 2012). Large-eddy simulations show that the LT homogenize the currents and leads to a negative bulk eddy viscosity near the surface (Sullivan et al., 2007) and suggests that LT has a nonlocal character. Several oceanic turbulence closure schemes incorporating LT
 45 have been proposed (e.g., McWilliams and Sullivan, 2000; Smyth et al., 2002; Harcourt and D’Asaro, 2008) and implemented into ocean models (Fan and Griffies, 2014;

Li et al., 2016). For K -profile parameterizations (KPP, see Large et al. 1994), a whole range of suggested modifications exist [see e.g., McWilliams and Sullivan 2000; Smyth et al. 2002; Harcourt and D’Asaro 2008; Takaya et al. 2010; Van Roekel et al. 2012],
 50 whereas Harcourt (2013) and Harcourt (2015) has introduced LT in a two-equation second-order turbulence closure schemes in the vein of Mellor and Yamada (1982).

1.2. *The wave-induced energy flux and turbulence*

The wave-induced turbulent kinetic energy (TKE) flux is usually parameterised in terms of a friction velocity (Craig and Banner, 1994). The parameterization works
 55 relatively well for fully developed sea state but performs poorly under growing and decaying sea states (Kantha and Clayson, 2004; He and Chen, 2011; Cai et al., 2017). The impact of the TKE flux is confined to the wave-breaking zone, and decays rapidly with depth (e.g. Terray et al., 1996; Janssen, 2012; Thomson et al., 2016; Esters et al., 2018). The importance of wave-generated turbulence near the sea surface was also
 60 demonstrated by Davies et al. (2000), for the bottom layer by Jones and Davies (1998) and for wave-induced turbulence by Babanin (2006), Huang and Qiao (2010), Babanin et al. (2010) and Babanin and Chalikov (2012).

1.3. *Stress balance at the air-sea interface*

The air-side wind stress is usually estimated from Monin-Obukhov Similarity Theory (MOST), e.g., COARE (Edson et al., 2013). When a wave model is available, the
 65 ocean surface roughness length can be estimated from the wave spectrum (Janssen, 1989). Previous studies (Drennan et al., 1999; Rutgersson et al., 2001; Höglström et al., 2013) have shown that the MOST may be invalid over waves. Some studies have tried to parametrize the total momentum flux over the ocean by two separate terms, i.e.,
 70 the wave-coherent momentum flux and the residual momentum flux (Höglström et al., 2015; Wu et al., 2017).

Traditionally, ocean circulation models assume that the air-side momentum flux (the wind stress τ_a) is identical to the momentum flux into the ocean interior, τ_{oc} . The role of surface waves is then ignored, assuming that surface waves do not take
 75 or release any momentum flux to ocean currents. This assumption is only valid when

surface waves are fully developed (i.e., in equilibrium with the local wind, Pierson and Moskowitz 1964). Mostly this is not the case, as winds are variable in time and space, and waves wax and wane on time scales from minutes to hours. Some of the momentum flux, τ_{in} is absorbed by the wave field, and so the net momentum flux to the currents in the ocean is (ECMWF, 2017; Staneva et al., 2017)

$$\tau_{oc} = \tau_a - \tau_{in} - \tau_{ds}, \quad (3)$$

where $\tau_a - \tau_{in}$ is the part of the wind stress τ_a that directly accelerates the ocean, and τ_{ds} is the momentum flux from the wave field to mean currents when the waves break. Here, τ_{ds} represents momentum loss from the wave field (Janssen, 2012).

The main effects of waves on the mean flow are due to radiation stress and Stokes drift, although interaction with turbulence and bottom stress can also be important (e.g., Wang and Shen, 2011). Surface waves can affect the water level and thus the storm surge through changes to the stress and upper ocean mixing and circulation (i.e., Dean and Dalrymple, 1991; Mastenbroek et al., 1993). Dietrich et al. (2011) showed that adding the radiation stress improves the model simulation of storm surges driven by hurricanes in the Gulf of Mexico, and a study of two wind storms in the North Sea (Staneva et al., 2017) found that adding the CSF and sea-state-dependent momentum and energy fluxes gave better agreement with measurements compared with a control experiment without wave effects.

Previous studies (Breivik et al., 2015; Alari et al., 2016; Staneva et al., 2017) have implemented the CSF and sea-state-dependent momentum and energy fluxes into the Nucleus of European Modelling of the Ocean (NEMO) model (Madec et al., 2015). Here, we add the Stokes drift in the mass and tracer advection equations with a layer-averaged Stokes drift profile (Breivik et al., 2016; Li et al., 2017). We will not attempt to investigate the impact of Langmuir turbulence in this study and will focus on the direct impact of the Stokes drift through the two competing processes of the CSF and material Stokes drift transport. This allows us to compare our results against these earlier studies which have left out the material Stokes drift transport. The wave-effects in the momentum equation (2) used in this study will be restricted to the Stokes-Coriolis

force. In the horizontal momentum equation, the vortex force and pressure correction
105 are of the order of the Rossby number relative to the Stokes-Coriolis force and so are
expected to be small on the spatial and temporal scales under consideration (see Section
2). We neglect the wave-effect in the vertical (hydrostatic) momentum equation stated
in Suzuki and Fox-Kemper (2016) in this study for simplicity. The rest of the paper
is structured as follows: the surface wave effects under consideration are presented
110 in Section 2. The experiment design as well as the measurements used in this study
are discussed in Section 3 and the simulation results are presented in Section 4. A
discussion of the most important findings is found in Section 5 and conclusions are
summarized in Section 6.

2. Wave effects introduced in the NEMO model

115 Four wave-related processes are investigated in the present study, (i) CSF, (ii) the
advection of tracer and mass by the Stokes drift, (iii) sea-state-dependent momentum
flux, and (iv) the flux of turbulent kinetic energy from breaking waves. In this study,
we are (a) focusing on the terms important on the larger scale and (b) not attempting to
investigate the detailed implications of Langmuir turbulence. We ignore the pressure
120 correction and vortex-force in the horizontal components of the momentum equation
(2) as they are of the order of the Rossby number. Although the model resolution
can become as fine as 3.7 km, we are not really resolving processes on that scale,
and certainly not the submesoscale (1–2 km) fronts where the Stokes-shear force has
been found to be important (Suzuki and Fox-Kemper, 2016; McWilliams and Fox-
125 Kemper, 2013; McWilliams et al., 2015). The Rossby number is generally small in our
model simulations so these terms may reasonably be neglected. Here we describe the
equations governing the evolution of NEMO as modified by the wave effects that we
include.

2.1. *Impact of Stokes drift on momentum and tracer and mass advection*

130 The Stokes drift impact on mean currents in large scale Eulerian ocean models is
through the CSF and LT. The CSF is a forcing induced by the interaction between the

Coriolis effect and the wave-induced Stokes drift (Hasselmann, 1970). The CSF can modify the vertical distribution of momentum in the turbulent Ekman layer (Polton, 2009). Previous studies (Breivik et al., 2015; Alari et al., 2016; Staneva et al., 2017) implemented the CSF in the NEMO ocean model but left out the additional effects of the Stokes drift on (i) the mass transport through the divergence of the sea-surface height, η , and (ii) the advection of tracers (e.g. McWilliams and Sullivan, 2000). After including the CSF (but not the other wave-effects), the wave-averaged momentum equation (2) simplifies to

$$\frac{D\mathbf{u}}{Dt} = -\frac{1}{\rho_w}\nabla p + \mathbf{u} \times f\hat{\mathbf{z}} + \mathbf{u}_S \times f\hat{\mathbf{z}} + \mathbf{D}^u. \quad (4)$$

The two additional effects (i) and (ii) of the Stokes drift in Eulerian ocean models can be written

$$\frac{Dc}{Dt} = -\mathbf{u}_s \cdot \nabla c + \mathbf{D}^c \quad (5)$$

$$\frac{\partial p}{\partial z} = -\rho g \quad (6)$$

$$\nabla \cdot \mathbf{u} = 0 \quad (7)$$

$$\nabla \cdot \mathbf{u}_s = 0 \quad (8)$$

$$\frac{\partial \eta}{\partial t} = -\nabla_h \int_{z=-H}^{z=\eta} (\mathbf{u}_h + \mathbf{u}_s) dz. \quad (9)$$

Here, c in the tracer advection equation (5) is a scalar quantity, such as temperature or salinity; and $\mathbf{D}^c$ is the parameterization of sub-grid scale physical processes.

The surface Stokes drift is usually parameterized using wind or friction velocity. Compared with the surface Stokes drift estimated from the full two-dimensional (2D) wave spectrum, the wind or friction velocity-dependent parameterizations overestimate or underestimate the surface Stokes drift under some wind conditions (see Appendix). The shear of the Stokes drift, an important factor when calculating the CSF, can be

calculated from the wave spectrum following Kenyon (1969). However, it is expensive
 150 to estimate the Stokes drift from the 2D wave spectrum as it has to be calculated at
 every vertical level of interest. Moreover, the full 2D spectrum may not be available.
 A monochromatic profile of the Stokes drift has commonly been used to estimate the
 Stokes drift profile in models. However, the shear of the monochromatic Stokes drift
 profile is much weaker than from a broad spectrum since the Stokes drift from short
 155 waves decays quickly with depth (Breivik et al., 2014). In order to estimate the Stokes
 drift more accurately, Breivik et al. (2016) proposed a more exact parameterization of
 the Stokes drift profile under a one-dimensional Phillips-type spectrum (Phillips, 1958)
 which has a better representation of the Stokes drift shear than the monochromatic pro-
 file (Breivik et al., 2016; Li et al., 2017). The Stokes drift profile is then approximated
 160 by,

$$u_S(z) = u_{s0} \left[e^{-2\bar{k}|z|} - \beta \sqrt{2\pi\bar{k}|z|} \operatorname{erfc}(\sqrt{2\bar{k}|z|}) \right], \quad (10)$$

where $\bar{k}$ is the inverse depth scale. The expression for $\bar{k}$ can be found in the study
 by Breivik et al. (2016). The parameter β is treated as constant equal to 1, following
 Breivik et al. (2016). The new Stokes drift profile has a better agreement with the pro-
 file calculated from a full wave spectrum compared than the monochromatic profile.
 165 As ocean models represent vertical and horizontal averages over a grid box, it is nec-
 essary to compute the layer-averaged Stokes drift profile, following Li et al. (2017),
 Appendix A:

$$u_S(z_i) = \frac{\int_{z_{i,b}}^0 u_S(z') dz' - \int_{z_{i,t}}^0 u_S(z') dz'}{z_{i,t} - z_{i,b}}. \quad (11)$$

Here $z_{i,b}$ and $z_{i,t}$ are the of the bottom and top of the i -th vertical level.

2.2. The sea-state dependent momentum flux

170 At the air-sea interface, the sea state modifies both the air-side momentum flux τ_a
 (by changing sea surface roughness and hence the drag coefficient) and the water-side
 momentum flux τ_{oc} (through the various amount of breaking waves under wave growth
 and decay). These two effects are described in the following subsections.

2.2.1. Water-side momentum flux

175 The momentum flux from the surface wave field to the ocean interior can be expressed in terms of the spectral source terms as

$$\tau_{oc} = \tau_a - \rho_w g \int_0^{2\pi} \int_0^\infty \frac{\mathbf{k}}{\omega} (S_{in} + S_{ds}) d\omega d\theta, \quad (12)$$

where S_{ds} is the dissipation due to surface wave breaking and depth-induced breaking, and S_{in} represents the wind input source term.

Fig. 1a illustrates the relationship between the wave age and the normalized water-side momentum flux, $|\tau_{oc}|/|\tau_a|$. The wind speed within the blue domain shown in Fig. 2 is shown as colored dots. The bin average with standard deviation is also shown in the figure as circles and bars, respectively. The normalized momentum flux ranges from about 0.85 to more than 1.8, which means that using the air-side momentum flux as a proxy will overestimate (underestimate) the water-side momentum flux by up to 20% (45%), respectively. In general, the ratio $|\tau_{oc}|/|\tau_a|$ decreases with increasing u^*/C_p , where u^* is the air-side friction velocity and C_p is the phase speed at the peak of the spectrum. Note that the standard deviation is larger for high values of u^*/C_p as there are fewer data points in this regime. Higher values of $|\tau_{oc}|/|\tau_a|$ occur mainly under low wind conditions.

190 2.2.2. Air-side momentum flux

The air-side momentum flux (wind stress) is usually calculated using the bulk formulae, i.e., $\tau_a = \rho_a C_D U_{10}^2$, where C_D is the drag coefficient. Following the MOST, the drag coefficient under neutral atmospheric conditions can be expressed as

$$C_D = \frac{\kappa^2}{\ln^2(z_{10}/z_0)}, \quad (13)$$

where κ is von Kármán's constant, $z_{10} = 10$ m and z_0 is the air-side surface roughness length. The surface roughness length varies with sea state. The Charnock relationship 195 (Charnock, 1955) is traditionally used to calculate the sea surface roughness length in

atmospheric models,

$$z_0 = \alpha \frac{u^{*2}}{g}, \quad (14)$$

where α is the Charnock coefficient. Here α is usually treated as a constant (0.018 is a common assumption, see Hersbach 2011). However, experiments show that it varies, usually in the range from 0.015 to 0.035 (Powell et al., 2003). The reason is that the sea surface roughness varies not just with wind speed but quite strongly with sea state (Janssen, 1989, 1991). Thus, many studies have parameterised the Charnock coefficient as a function of sea state, such as wave age (e.g. Guan and Xie, 2004). In the coupled ECWAM model (ECMWF, 2017), the Charnock coefficient is parameterised according to quasi-linear theory as a function of the air-side momentum flux that goes into the wave field,

$$\alpha = \frac{\hat{\alpha}}{(1 - \tau_{\text{in}}/\tau_{\text{a}})^{1/2}} \quad (15)$$

where $\hat{\alpha} = 0.006$, and the wave-induced stress is expressed as (Janssen, 1991),

$$\tau_{\text{in}} = \rho_{\text{w}} g \int_0^{2\pi} \int_0^\infty \frac{\mathbf{k}}{\omega} S_{\text{in}} d\omega d\theta \quad (16)$$

where ρ_{w} is the water density, ω the angular frequency, θ the wave direction, and S_{in} is the wind input source term. The air-side stress can now be found by iteratively by calculating the roughness length (14) from the modified Charnock parameter (15). See Janssen (2004), pp 122–124 for further details on the coupling of ECWAM to the atmospheric model. In our runs the wave model is uncoupled from the atmospheric model, therefore no feedback to the atmospheric boundary layer takes place, but the air-side stress τ_{a} is calculated using the sea-state dependent roughness length following Eqs. (13)–(16).

An example of the relation between the wave age and α is shown in Fig. 1b. The hourly output data from the German Weather Service (Deutscher Wetterdienst, DWD) forcing during December 7th, 2015 in the blue box domain shown in Fig. 2 are indicated as dots (color represents wind speed, U_{10}). Generally, α increases with increasing inverse wave age (u^*/C_{p}) as the sea surface grows rougher when the sea state becomes

younger. However, the α has significant scatter under the same wave age caused by differences in wave steepness (the standard deviation and the bin average of α are shown as circles and bars in Fig. 1b). Thus, parameterizations based on bulk estimates will give the same value even if the wave spectra differ significantly. Some studies (e.g.,
225 Edson et al., 2013) show that the Charnock coefficient asymptotes to 0.03 at high wind speed. However, α from the WAM model is larger which may be due to that the sea spray influences are not included when calculating α in the WAM model.

2.3. Sea-state-dependent energy flux

The TKE flux is ejected into ocean currents due to breaking waves (Terray et al.,
230 1996; Drennan et al., 1996; Sutherland and Melville, 2015). Craig and Banner (1994) suggests that the TKE flux can be parameterized using the air-side friction velocity as $m\rho_a u_*^3$. If it is assumed that the water side is in balance with the air side (waves in equilibrium with the local wind), the coefficient m is usually treated as constant with a value of about 3.5. Likewise, the water-side friction velocity, u_* , can be used, in which
235 case the flux is written as $\alpha_{CB}\rho_w u_*^3$. Since the momentum flux across the surface is continuous, i.e., $\rho_a u_*^2 = \rho_w u_*^2$, and it is clear that $\alpha_{CB} = \sqrt{(\rho_w/\rho_a)}m \approx 100$ if assuming $m = 3.5$ (Craig and Banner, 1994). Studies have shown that the coefficient α_{CB} depends on the sea state (Terray et al., 1996; Feddersen et al., 2007; Jones and Monismith, 2008; Gerbi et al., 2009; Fan and Yu, 2017) and is in the range ≈ 70 to
240 over 200. The coefficient α_{CB} can be calculated from the spectral source terms in a wave model (Janssen, 2012; Breivik et al., 2015),

$$\Phi_{oc} = -\rho_w g \int_0^{2\pi} \int_0^\infty S_{ds} d\omega d\theta = -\alpha_{CB}\rho_w u_*^3 \quad (17)$$

The variability of α_{CB} for the North Sea and the Baltic Sea study areas has been discussed by Staneva et al. (2017) and Alari et al. (2016).

3. Experiments and data

3.1. Model setup

245 NEMO is a state-of-the-art framework of ocean-related model components. The OPA (physical core engine) and LIM3 (sea-ice dynamics and thermodynamics pack-

age) packages are used in this study to investigate wave effects on ocean circulation and hydrography in a regional high-resolution model domain. The details of the NEMO model are described by Madec et al. (2015) and the details of the North Sea-Baltic Sea set-up are documented by Staneva et al. (2017). The domain used in this study is shown in Fig. 2, which covers the Baltic Sea and North Sea. The horizontal resolution of the model is 2 nautical miles (about 3.7 km) with 56 z -levels in the vertical. The vertical layer thickness ranges from 1.5 m at the surface to around 22 m in the lowest model levels with 5 levels in the top 10 m. The temperature and salinity climatology by Janssen et al. (1999) is used as initial conditions. At the open boundaries, the Oregon State University (OSU) tidal constituent data base is used (Egbert and Erofeeva, 2002) as well as climatological periodic boundary conditions for the temperature and salinity (Janssen et al., 1999). Hourly atmospheric forcing data (i.e. wind fields, shortwave and long-wave radiation, air temperature, humidity and air pressure) are from DWD’s short range regional COSMO-EU atmospheric forecasts with data assimilation (Baldauf et al., 2011). COSMO-EU has a 7 km horizontal resolution and 40 vertical levels up to about 24 km with a domain covering the whole of Europe. COSMO-EU is nested into a global model with 30 km grid resolution¹. A $k - \epsilon$ vertical diffusion scheme (prognostic equations for TKE, k , and turbulent dissipation, ϵ) is used for all simulations through the generalized length scale (GLS) scheme (Umlauf and Burchard, 2003, 2005).

The Wave Model (WAM) (WAMDI, 1988; Komen et al., 1994; Janssen, 2004) is a third-generation spectral wave model. The 2D wave spectrum evolves according to the energy balance equation,

$$\frac{dF}{dt} + \frac{\partial}{\partial \phi}(\phi F) + \frac{\partial}{\partial \lambda}(\lambda F) + \frac{\partial}{\partial \theta}(\theta F) = S_{\text{in}} + S_{\text{ds}} + S_{\text{nl}} \quad (18)$$

where F represents the wave spectral density, ϕ denotes the latitude, λ represents the longitude, and S_{nl} is the nonlinear transfer term. The wind input and dissipation terms

¹For more details please refer to https://www.dwd.de/EN/research/weatherforecasting/num_modelling/01_num_weather_prediction_models/regional_model_cosmo_eu.html

in the WAM model are described in Janssen (1989, 1991). The nonlinear term is based on the parameterization proposed by Hasselmann et al. (1985).

275 The domain and resolution of the WAM model are the same as for NEMO with 24 directions and 25 frequencies for the wave spectrum (Staneva et al., 2017). The 10 m wind vectors from DWD’s short range forecasts are used to force WAM. The required boundary information for the wave model at the open boundaries of the North Sea is taken from the hourly output of the regional wave model EWAM (Europe WAM),
280 which is run twice a day in operational wave forecast routine at DWD. Hourly fields of surface Stokes drift, significant wave height and mean frequency yield Stokes drift profiles while fluxes of momentum and TKE are used to force NEMO and to increase mixing under breaking waves (shown in Section 2).

3.2. Experiments

285 Four experiments were devised to investigate the wave impact on NEMO (see Table 1). The CTL is the stand-alone NEMO run without explicit wave information. The bulk algorithm developed by Large and Yeager (2004, 2009) is used to calculate the surface fluxes. The normalized energy flux parameter from breaking waves, α_{CB} , is treated as 100 in CTL, STCOR and STFUL (Craig and Banner, 1994). The CSF
290 is introduced in STCOR (as expressed in Eq. 2) but no mass or tracer advection by the Stokes drift. In the STFUL experiment, the Stokes drift is introduced in the mass and tracer advection (equations 9 and 5) in addition to the CSF. All wave processes discussed in Section 2 are implemented in experiment FULL.

The default model (CTL) was used to simulate the period 2014-01-01 to 2015-12-
295 31. To investigate the wave effects, the following periods were run:

- Storm period simulation (2015-10-31 to 2015-11-30):

The wave effects were switched on from the restart file of the CTL experiment at time 2015-10-31 00:00. The three experiments, i.e, STCOR, STFUL and FULL, are here used to investigate wave effects under storm conditions. During this
300 time period, there are a few low pressure systems passing through the North Sea and the Baltic Sea. Fig. 3(a) and (b) shows the time series of the wind speed and significant wave height at the FINO1 station. The maximum wind speed

exceeds 20 m s^{-1} while the significant wave height reaches 7 m. During this time period, the inverse wave age u^*/C_p is between 0 and 0.1. The larger value of u^*/C_p is mainly in the low wind conditions (Fig. 3(c)). The larger value of α_{CB} and u_{s0} corresponds to the young waves and high wind conditions (Fig. 3(d) and (f)). For the normalized wind stress, τ_{oc}/τ_a , the highest values appear for the oldest waves (lowest inverse wave age) and conditions with sudden wind changes (Fig. 3(e)).

• Seasonal simulation (2015-01-01 to 2015-12-31):

The wave effects were switched on at 2015-01-01 00:00 for the seasonal simulations. The three experiments, i.e., STCOR, STFUL and FULL, and the control simulation (CTL) during year 2015 were used to analyze the seasonal impact of waves on circulation and the Baltic Sea upwelling. The seasonal mean fields of U_{10} , significant wave height and swell height are shown in Fig. 4. During months JFD (January, February, and December), the mean wind speed in the North Sea is higher than 12 m s^{-1} and less than 10 m s^{-1} in the Baltic Sea. The mean wind direction is from southwest to south in the North and Baltic Seas. The significant wave height (H_s) decreases from the north (more than 3 m) to the south (less than 2 m) in the North Sea. This may be due to swell propagating from the Norwegian Sea to the North Sea (Fig. 4c shows the swell significant wave height). The swell significant wave height, H_{swell} , is higher than 1.6 m in the north part of the North Sea, and less than 0.8 m in the Baltic Sea. In MAM (March, April and May), the mean wind speed is less than 9 m s^{-1} , and the direction is from southwest both in the North and Baltic Sea. The mean H_s is less than 1.8 (1.4) m in the North Sea (Baltic Sea). In JJA (June, July and August), the mean direction is from the northwest and the wind speed is reduced to less than 7 m s^{-1} . The H_{swell} in the Baltic Sea is less than 0.7 m. In SON (September, October and November), the mean wind is from the southwest and the wave height for both swell and total wave height is larger than in JJA.

3.3. Data

The following measurements were used to assess the model performance in this study.

- MARNET

335 The Marine Environmental Monitoring Network in the North Sea and Baltic Sea (MARNET) comprises twelve automated measuring stations. It is operated by the Federal Maritime and Hydrographic Agency of Germany (BSH). The station locations are shown as red dots in Fig. 2. Their locations are listed in Table 2. The data include water temperature, salinity, oxygen saturation, and meteorological data. For details of the stations and data refer to http://www.bsh.de/en/Marine_data/Observations/MARNET_monitoring_network/index.jsp.
340

- ADCP

Measurements from Acoustic Doppler Current Profilers (ADCP) at FINO3 were
345 used to compare with the current velocities in NEMO. The type of ADCP used (Nortek Acoustic Wave and Current Meter (AWAC)) is designed to measure both current profiles and wave parameters.

- SMHI dataset

Hourly sea level measurements from the dataset of the Swedish Meteorological
350 and Hydrological Institute (SMHI) were also used. The sea level data sites are marked as blue dots in Fig. 2. Their locations are listed in Table 2. For details see <https://www.smhi.se/en/services/open-data/oceanographic-observations-1.33356>.

4. Results

355 4.1. Extreme conditions

Fig. 5 shows the sea level anomalies (SLA) at station Furuogrund during the storm simulation period. In general, the CTL experiment reproduces the SLA pattern well. During the storms, CTL underestimates the SLA by up to 30 cm at this station. Adding

the CSF influence (STCOR) lowers the SLA (increases the difference with observations
 360 compared with the results from the control experiment). Adding both the CSF and the
 Stokes drift advection (STFUL) yields only small differences compared with the CTL
 experiment in terms of SLA because the impact of the Stokes advection is to counter
 the impact of the CSF. Adding the sea-state-dependent momentum and energy fluxes
 to STFUL in experiment FULL reduces the SLA bias by up to 20 cm during the storm
 365 period at station Furuogrund. As demonstrated by Staneva et al. (2017), the sea-state-
 dependent momentum flux has a larger impact than the sea-state-dependent TKE flux.

The comparison between the four model experiments and the observations at the
 24 stations (shown as blue dots in Fig. 2) in the Baltic Sea during the integration period
 from 2015-10-31 to 2015-11-30 are shown in a Taylor diagram (Fig. 6). The Taylor
 370 diagram quantifies three statistics (i.e., Pearson correlation coefficient, the centered
 root mean square error, and the standard deviation) between the modelled and observed
 results. In general, adding the CSF (STCOR, represented by orange) has some impact
 on the model performance concerning standard deviation, centered root mean square
 error and correlation compared with the control run (blue), but there is no general
 375 improvement. In some stations, adding the CSF improves the model performance but
 it worsens the results in other stations (see Fig. 6). Adding both the influence of CSF
 and the Stokes advection (STFUL, yellow) slightly improves the model performance
 in terms of SLA. The model performs best when the sea-state-dependent momentum
 and energy fluxes are added (FULL, purple). FULL reduces the standard deviation and
 380 root mean square error significantly. However, none of the three experiments (STCOR,
 STFUL and FULL) increase the correlation significantly.

The velocity profiles at FINO3 during days 315-320 are shown in Fig. 7. One can
 see that all the experiments catch the changing velocity pattern during the simulation
 period. During day 318, the currents vary strongly with depth (Fig. 7a), which is not
 385 captured by any of the four experiments. In general, the control experiment (Fig. 7b)
 overestimates (underestimates) the current velocity at the high (low) velocity peaks at
 FINO3. Adding the CSF does not have a systematic impact on the simulated currents.
 For example, adding the CSF (STCOR, Fig. 7c) increases the current velocity during
 the high-velocity periods (the warm colour period) for days 317–319, but decreases

390 slightly the currents at day 317.3 and 318. In general, adding both CSF and the Stokes
 advection (STFUL) reduces the impact of the CSF. However, the impact of the Stokes
 drift on the mass and tracer advection does not cancel the influence from the CSF
 (comparing the simulation results from CTL and STFUL). Adding all the wave effects
 (FULL) generally leads to better agreement with the measurements compared with the
 395 experiments STCOR and STFUL.

The net heat flux and the mixed layer depth at 01:00 UTC, day 318 (November
 14), is shown in Figs. 8 and 9. The spatial patterns seen here are representative of the
 full storm period. The CSF decreases slightly the heat flux in the centre of the North
 Sea and the south coast of Sweden (Fig. 8b) whereas when both the CSF and Stokes
 400 advection are added (STFUL), the impact is negligible (Fig. 8c). In general, adding all
 the wave effects (FULL) decreases the heat flux, especially in the North Sea and the
 southern part of the Baltic sea (by up to 40 W m^{-2} , see Fig. 8d). Due to the changes of
 the heat flux and the momentum flux, adding all the wave influences (FULL) increases
 the mixed layer depth by more than 10 m (more than 10%) in the North Sea (Figs.
 405 9d). The mixed layer depth is defined as the layer where the potential temperature is
 $0.01 \text{ }^{\circ}\text{C}$ different from the SST. However, the CSF and Stokes advection do not have a
 significant impact on the mixed layer depth (Figs. 9b and c). The sea-state-dependent
 momentum flux changes the stress and indirectly changes the heat flux and sea temper-
 ature. As stated by Breivik et al. (2015), the ocean model has too vigorous mixing with
 410 a constant α_{CB} . Compared with the wind stress, sea-state-dependent TKE flux has less
 impact on SST since its influence is mainly on the surface layer. The influence of the
 sea-state-dependent momentum flux and TKE flux varies geographically. In general,
 compared with sea-state-dependent TKE flux, the sea-state-dependent momentum flux
 has a greater impact on both the heat flux and ocean mixed layer depth, which agrees
 415 with the results of Alari et al. (2016).

As shown by Alari et al. (2016), the surface wave influence can extend to the bottom
 layer, but the main impact is in the mixed layer. Here, the relative mean difference
 (RMD) of SST and MLD between the control run and FULL are used to show the
 wave impact with different wave parameters (see Fig. 10). Hourly model fields from
 420 within the domain shown as a blue box in Fig. 2 are used. One can see that the wave

impact on the SST decreases with the mean square slope (σ^2) in the range [0 0.01] (from $\approx 2.8\%$ to 1%), then it increases with σ^2 until $\sigma^2 = 0.05$. A similar pattern is found for RMD of MLD (see Fig. 10b). In general, the RMD of MLD and SST increases with the increasing of the value $|\tau_{oc}/\tau_a - 1|$ (Figs. 10c and d). This is due to the ocean-side stress is not equal to the air-side stress which differs the energy input to the ocean compared with the control run. The ratio H_{swell}/H_{sea} shows the relative energy of swell waves. It shows that the RMD of SST and MLD decreases with increasing of the relative swell energy (see Fig. 10e and f). The wave impact on the ocean simulation in terms of SST and MLD first increases with inverse wave age (u^*/C_p) and then decreases with inverse wave age. The maximum influence occurs at $u^*/C_p = 0.06$.

4.2. Seasonal influence

The 3-hourly averaged output from the seasonal simulation period is used to investigate the wave effects. Fig. 11 shows the SST discrepancies between the model and measurements for all stations marked with red circles in Fig. 2. The correlation coefficients for the four experiments are all higher than 0.9. There is no significant difference between the experiments STCOR/STFUL and CTL concerning the root mean square error of SST, standard deviation and coefficient at those stations. Adding the sea-state-dependent momentum and energy fluxes improves the model simulation of SST (increases the correlation coefficients and decreases the standard deviation and root mean square errors).

The mean summer SST (JJA) is shown in Fig. 12a. In summer, the SST along the Swedish south coast is lower than the zonal average SST in the Baltic Sea. This may be caused by the coastal upwelling. Adding the CSF (STCOR) increases the SST by more than $0.2\text{ }^\circ\text{C}$ in the southern Baltic Sea and in some areas in the western North Sea. In contrast, adding the CSF reduces the SST along the Finnish coast and the south coast of Norway (Fig. 12b). Adding both the CSF and the Stokes advection, STFUL, reduces the influence from the CSF (Fig. 12c). The SST differences between FULL and CTL show no significant difference (less than $0.05\text{ }^\circ\text{C}$). When the sea-state-dependent momentum and TKE fluxes are added to STFUL (FULL), the SST

is reduced by more than 0.6 °C in the Baltic Sea and 0.3 °C in the North Sea, away from the coast. However, the FULL experiment increases the SST along the eastern and southern Swedish coast, the east coast of Gotland, and the Finnish south coast in summer (Fig. 12d).

455 Fig. 13a shows the mean SST from the CTL experiment in winter (JFD). Adding the CSF increases the SST about 0.3 °C in the southern and northern parts of the Baltic sea, but decreases the SST by about 0.2 °C in the middle part of the Baltic Sea (see Fig. 13b). Similar to the summertime period, adding the Stokes advection reduces the influence of CSF on the SST (Fig. 13c). Adding all the wave effects (FULL) increases
460 the SST more than 0.4 °C in the Baltic Sea. However, adding all the wave effects increases the SST in some areas in the North Sea (e.g., the south part of the North Sea and the southern coast of Norway) and decreases the SST in some areas in the North Sea (see Fig. 13d).

Comparing the SST from the three wave experiments (STCOR, STFUL, and FULL)
465 with the CTL experiment, we see that as expected the Stokes advection reduces the influence of the CSF. The sea-state-dependent momentum and TKE fluxes dominate over the CSF and the Stokes advection in this study.

4.3. Coastal Upwelling

Coastal upwelling is an important phenomenon. Where the wind blows parallel
470 to a coastline, wind-driven horizontal divergence can cause cold, nutrient-rich water to rise to the surface. Upwelling can lead to a drop in the sea surface temperature (SST) of more than 10°C in less than two days (Lips and Lips, 2008). The divergence (convergence) of the wind stress is the major factor affecting upwelling (Lehmann and Myrberg, 2008). Upwelling can affect the marine atmospheric boundary layer (Spro-
475 son and Sahlée, 2014), CO₂ fluxes (Norman et al., 2013) and biological productivity (Lehmann and Myrberg, 2008). The Baltic Sea is a semi-enclosed basin and winds from all directions cause upwelling along one coast or another. Climatological studies have shown that the Baltic Sea has high upwelling frequencies in the summer months (e.g., Lehmann and Myrberg, 2008). The surface wave influence on the upwelling has
480 rarely been investigated in previous studies. Here, we have focused on the impact of

surface waves on the upwelling events which have a surface temperature signal in the Baltic Sea. The method described in Lehmann et al. (2012) was used to detect upwelling events. A north-south SST gradient is usually present in the Baltic Sea. We use the SST difference in every model grid point from the zonal mean to detect upwelling events. Different SST criteria have been used in other studies, such as 1 °C (Norman et al., 2013). Here, we require that the SST difference be greater than 2.5 °C within 25 km from the coast.

With 3-hourly resolution, we record upwelling events in all model grid points within the Baltic Sea. Then the upwelling frequency is estimated as the ratio between the number of upwelling events and the total number of samples in each grid point. The upwelling frequency from the control experiment (CTL) for the months June, July, August and September 2015 is shown in the first column of Fig. 14. In June and July 2015, the most frequent upwelling is found along the Swedish south coast, the east and south coasts of Öland and Gotland islands. The upwelling frequency is generally above 30%, with extremes of 80% in those areas during June and July 2015. In June 2015, the upwelling frequency exceeds 50% along the coast of Helsinki. In August 2015, the upwelling frequency along the Swedish south coast is below 40%. In September 2015, upwelling events are mainly along the east coast of Öland, with frequency in excess of 50%, and along the coast of Gotland with a frequency a little below 30%. In August and September, the upwelling events in the Gulf of Finland are quite infrequent. One possible reason is that the upwelling events are weak in those two months and cannot be detected using 2.5 °C temperature difference with zonal mean as the upwelling criterion.

The second to fourth columns of Fig. 14 show the wave impact on the upwelling frequency for the months June, July, August and September 2015. Warm colours indicate that adding the wave effect increases the upwelling frequency and cold colours indicate that it reduces the upwelling frequency. In June, adding the CSF (Fig. 14b) reduces the upwelling frequency (about 3% and extremes of around 5%) along the Swedish south coast. In July, it reduces the upwelling frequency by more than 5% and in extreme cases by about 10% (Fig. 14f). Adding the CSF increases the upwelling frequency in September along the east coast of Öland. In the Gulf of Bothnia and Finland

and along the coast of Gotland, STCOR generally increases the upwelling frequency compared with CTL. The maximum frequency difference induced by the CSF occurs in September, at most by about 6% in the Gotland coast. Adding both the CSF and the Stokes drift advection (STFUL) generally reduces the impact of CSF, but there are some exceptions in the Gulf of Finland and along the south coast of the Gotland. In the Gulf of Finland, adding both the CSF and the Stokes drift impact on the mass and tracer advection has a bigger impact on the upwelling frequency than adding only the CSF. Along the south coast of the Gotland, it even reverses the trend from increasing the upwelling frequency (STCOR) to decreasing it (STFUL). When all the wave effects are added (FULL), a bigger influence on the upwelling frequency is found in June, July and August. In June, the experiment FULL reduces the upwelling frequency along the south Swedish coast by more than 8%. FULL actually increases the upwelling frequency compared with CTL (see Fig. 14h) along the Swedish east coast in July. In September, adding only the CSF has a greater impact on the upwelling frequency than adding both the CSF and Stokes drift tracer advection (STFUL) or adding all wave effects (FULL), see Fig. 14n, o, and p).

In general, adding CSF changes the upwelling frequency significantly in the areas that surface Stokes drift gradient is large, such as the east and south coast of Sweden in June, July and August (see the first column of Fig. 15). The Stokes drift in the Swedish coast is smaller than the areas far away from the coast. It indicates that the CSF-induced upper ocean mixing is smaller in the coastal areas than the others. In most cases (see also Figs. 12c, 13c), the CSF and the Stokes drift material and tracer advection tend to counter each other. The sea-state dependent TKE and momentum fluxes dominate wave-related processes investigated in this study. One can see that the momentum and energy fluxes are reduced in the coastal areas when adding sea-state dependent momentum/TKE flux. This leads to less upwelling and increases the upper ocean mixing in the areas far away from the coast. Thus, they reduce the upwelling frequency (see the second and third columns of Fig. 15). It is worth noting that the sea-state dependent stress increases the stress curl in the areas (the upwelling detected areas, within 25 km from the coast) where τ_{oc}/τ_a changes quickly. This drives the increase in the upwelling frequency along the Swedish east coast in July. It indicates that

younger waves along the coast than other areas can potentially decrease the upwelling frequency. However, the stress curl induced by the waves can also contribute to the upwelling frequency difference. As the sea state modifies both the spatial distribution of wind stress and the ocean interior through the CSF, the combined effect is hard to predict and may vary over short distances due to topographic features.

The SST difference from the zonal mean temperature (ΔT) is used as an index to assess the wave impact on the upwelling intensity in the Baltic Sea. Here, we divide the upwelling events into three different ranges according to the SST difference: weak upwelling ($-4 < \Delta T < 2.5^\circ\text{C}$), intermediate upwelling ($-6 < \Delta T < -4^\circ\text{C}$), and strong upwelling ($\Delta T < -6^\circ\text{C}$). Fig. 16 shows the normalised distribution of the upwelling events in the months June, July, August and September 2015. More than 60% of upwelling events are weak in these four months. Comparing with the months July, August and September, there are more weak upwelling events (about 78% of the total upwelling events) and fewer strong upwelling events (about 1.2%) in June 2015. In July, there are more intermediate upwelling events than in other months, which is about 31.5% of the total upwelling events (it is less than 27% in the other three months). The strong upwelling events represent about 7% of the total in September, but it is only 1.7% in June and 4.6% in August.

The impact of waves on the upwelling intensity distribution in the months June, July, August and September are shown in Fig. 17. Compared with CTL, adding the CSF (STCOR) increases (less than 1%) the weak upwelling probability in June, July and September, but decreases slightly the weak upwelling in July. In general, the CSF reduces the intermediate and strong upwelling probability but increases by about 1% the strong upwelling in August. Adding both the CSF and the Stokes drift advection (STFUL) changes the upwelling only slightly. This agrees with the results from previous sections where it was shown that the Stokes drift advection counters the influence of the CSF. When all the wave effects are added (FULL) a greater impact is found. FULL increases the weak upwelling probability by about 9% in July and 8% in August and September, but only by about 5% in June. It reduces the intermediate (strong) upwelling probability by about 4% in June, 6% (2%) in July, and 5% (1%) in August and September.

5. Discussion

575 As one of the most important wave-related processes, the CSF can modify the ocean Ekman transport and lead to an additional veering of the Ekman profile by about $4\text{--}5^\circ$ (McWilliams and Restrepo, 1999; Polton et al., 2005; Sætra et al., 2007). The current profile can thus be substantially modified in the whole wind-driven layer as shown by Polton et al. (2005). This is also evident in Fig. 7. However, due to indirect
580 influences from CSF, the CSF impact on the current in the location shown in Fig. 7 are not significant (but there are still some differences). The influence varies with location. The implication is that the impact of Stokes drift on mixing and the Eulerian current via Langmuir turbulence may dominate over this effect. In addition, the Stokes drift may indirectly affect the mixing by modifying the Eulerian shear. This is quite
585 evident in Figs. 12b and 13b where we see that the inclusion of the CSF modifies the SST throughout the entire North Sea and the Baltic Sea and not just in the expected upwelling and downwelling areas illustrated in Fig. 14.

The influence of the Stokes drift varies with location and time. The surface Stokes drift is mainly controlled by the high-frequency part of the wave spectrum. Thus, the
590 directional difference between the wind and the Stokes drift is smaller under growing waves since the local wind waves dominate the Stokes drift. When swell dominates, the Stokes drift contribution by low-frequency waves can be substantial, and their direction will in general not align with the local wind direction which means a simple wind-speed parameterization of the Stokes drift will tend to perform poorly. In other words,
595 the wave model may yield a bigger influence under swell conditions where wind and waves are not aligned.

In general, adding the wave effects has a positive impact on the model performance compared with the control experiment. However, the biases between the model results and measurements are still large. The biases may be due to the setup of the NEMO
600 model, the bulk parameterization or the forcing data. In this study, our aim was not to achieve the best fit between data and models but rather to investigate the role of the wave-induced processes on hydrography, circulation and its impact during storm events. In addition, we have not included Langmuir turbulence and bottom stress in-

duced by wave-current interaction in shallow water (Davies and Lawrence, 1995), nor
 605 have we considered the possible impact of non-breaking waves put forward by Qiao
 et al. (2004), Babanin (2006) and Babanin and Haus (2009). Wave effects can affect
 each other, making wave-current interaction a complex non-linear feedback system.
 For example, due to the Langmuir turbulence, the turbulent mixing of Eulerian momen-
 tum occurs down both the Eulerian current and the Stokes drift gradient (McWilliams
 610 et al., 2014; Harcourt, 2015; Reichl et al., 2016). The down-Stokes gradient mixing
 may partially cancel the Stokes advection because it introduces a counter-Stokes effect
 on the Eulerian current which is because that mixing acts to homogenize the combined
 Lagrangian current. Further studies are needed to disentangle the full impact of these
 wave-related processes on ocean circulation.

615 **6. Conclusions**

This study extends the work of Alari et al. (2016) and Staneva et al. (2017) on a
 regional NEMO model with WAM forcing by also introducing the impact of Stokes
 drift on the mass and tracer advection. Using the NEMO model with a domain cover-
 ing the Baltic Sea and North Sea we have investigated the impact of the CSF, Stokes
 620 drift on the mass and tracer advection and sea-state-dependent momentum and energy
 fluxes on ocean circulation on very high resolution. Of those wave-related processes,
 the sea-state-dependent fluxes dominate the Stokes drift effects. In more detail, the
 sea-state-dependent momentum flux has a larger influence on the simulation than that
 from sea-state-dependent TKE flux, which agrees with previous studies (Alari et al.,
 625 2016; Staneva et al., 2017). Furthermore, we found that overall the Stokes drift ad-
 vection largely counters the effect of the CSF. The CSF can change the SST by more
 than 0.2 °C, and the influence varies with location and season. The Stokes drift im-
 pact on the mass and tracer advection tends to counter the influence of the CSF on
 ocean circulation. However, these two effects do not always cancel entirely, as seen for
 630 the SST (Figs. 12c and 13c) and the associated coastal upwelling (Figs. 14 and 17).
 Thus, care must be taken to consistently introduce these wave effects together since
 their combined effect is much smaller (yet still present) than either on its own. The

terms neglected in this study (e.g., the Stokes correction to the pressure in the momentum equation and the wavy-hydrostatic effect) will be investigated in further studies.

635 Adding all the wave effects discussed in Section 2 decreases (increases) the mean SST by more than 0.6 (0.4) °C in the Baltic Sea during summer (winter).

Including CSF in NEMO changes the Baltic sea upwelling frequency by more than 10%. The combined impact of the CSF and the Stokes drift impact on mass and tracer advection reduces the upwelling frequency difference to less than 2% compared with

640 the control experiment. When all the wave effects are added, the upwelling frequency changes by more than 10% compared with the control experiment. In addition, the wave effects increase (reduce) the weak (strong) upwelling probability, which is mainly caused by the sea-state-dependent momentum and energy fluxes.

The sea-state-dependent momentum and energy fluxes dominate the wave-induced

645 processes investigated in this study. Introducing the sea state in the calculation of the momentum and energy fluxes paves the way for consistent coupled regional atmosphere-wave-ocean models. As it changes the divergence (convergence) of the wind stress this can modify the upwelling frequency by more than 10%.

In general, adding wave effects improves the model performance compared with the

650 control experiment. The CSF and the Stokes drift mass and tracer advection have only a small impact on the SLA during storms. The sea-state-dependent fluxes improve the modelled sea level. These results are consistent with the conclusions by Mastenbroek et al. (1993) and Saetra et al. (2007).

In this study, NEMO was forced by WAM, and there is no feedback from the ocean

655 interior (NEMO) to the wave field (WAM). In reality, the atmosphere and ocean form a coupled system with the wave field as the mediator. Surface waves affect the wind by changing the air-side roughness length (Janssen, 1989, 1991; Wahle et al., 2017). Changes to the upper ocean brought about by the presence of waves, in terms of temperature and near-surface currents, will, in turn, affect the atmospheric boundary layer

660 and ultimately the evolution of low and high-pressure systems (Sheldon and Czaja, 2014). Also, surface waves are affected by ocean currents through current refraction (Janssen, 2004). These processes can only be investigated with fully coupled systems (Mogensen et al., 2017).

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8. Appendix

The Stokes drift must be parameterised from the local wind when wave information
675 is not available. The magnitude of the surface Stokes drift, i.e., u_{s0} , is usually estimated
through a wind-speed dependent parameterization, e.g., $u_{s0} = 0.016|U_{10}|$ (Li and Gar-
rett, 1993), or a stress-dependent parameterization, e.g., $u_{s0} = 0.377|\tau/\rho_w|^{1/2}$ where
 $|\tau|$ is the surface wind stress (Madec et al., 2015) and ρ_w is the water density. The
direction of Stokes drift is then assumed to align with the wind direction. However,
680 those parameterizations are unable to capture the Stokes drift in the presence of swell
($C_p/(\cos(\theta_{wave} - \theta_{wind})U_{10}) > 1.2$, where C_p is the wave peak phase speed), as
the direction and magnitude of the Stokes drift can be altered significantly by swell
propagating in off-wind directions, say $|\theta_{wind} - \theta_{wave}| > 45^\circ$.

Fig. 18 shows the relationship between the magnitude of the surface Stokes drift
685 estimated by the WAM model and the mean wind speed (18a) and friction velocity
(18b). The dots represent the output from all grid points within the blue box shown
in Fig. 2. The wind forcing of WAM is from a short range forecast by the German
Weather Service (DWD) on December 7, 2015. In Fig. 18, the dots represent all
model values inside the blue domain shown in Fig. 2. The surface Stokes drift is
690 calculated from the full 2D wave spectrum in WAM. The contribution from the spectral
tail beyond the cut-off frequency to the surface Stokes drift is calculated diagnostically
and added to the contribution from the main part of the WAM spectrum (ECMWF,

2017). The color represents the wave age, i.e., C_p/U_{10} . The black lines represent the results from the parameterization $u_{s0} = 0.016U_{10}$ (Li and Garrett, 1993) (Fig. 18a) and $u_{s0} = 0.377|\tau/\rho_w|^{1/2}$ (Madec et al., 2015) (Fig. 18b). There is significant scatter between the results from the parameterizations and that calculated by the 2D wave spectrum. Especially under low wind conditions, both parameterizations overestimate the surface Stokes drift. The direction of the surface Stokes drift differs from the mean wind direction, especially under high wave age conditions (Fig. 18c). The reason is that swell propagating in a different direction than the local wind contributes to the Stokes drift. Thus, the surface Stokes drift parameterised either from the mean wind speed or from the friction velocity can not properly model the Stokes drift.

References

- Alari, V., Staneva, J., Breivik, Ø., Bidlot, J.R., Mogensen, K., Janssen, P., 2016. Surface wave effects on water temperature in the baltic sea: simulations with the coupled nemo-wam model. *Ocean Dynamics* 66, 917–930.
- Axell, L.B., 2002. Wind-driven internal waves and langmuir circulations in a numerical ocean model of the southern baltic sea. *Journal of Geophysical Research: Oceans* 107.
- Babanin, A., 2006. On a wave-induced turbulence and a wave-mixed upper ocean layer. *Geophysical Research Letters* 33.
- Babanin, A.V., Chalikov, D., 2012. Numerical investigation of turbulence generation in non-breaking potential waves. *Journal of Geophysical Research: Oceans* 117.
- Babanin, A.V., Chalikov, D., Young, I., Savelyev, I., 2010. Numerical and laboratory investigation of breaking of steep two-dimensional waves in deep water. *Journal of Fluid Mechanics* 644, 433–463.
- Babanin, A.V., Haus, B.K., 2009. On the existence of water turbulence induced by nonbreaking surface waves. *Journal of Physical Oceanography* 39, 2675–2679.

- Baldauf, M., Seifert, A., Förstner, J., Majewski, D., Raschendorfer, M., Reinhardt, T.,
720 2011. Operational convective-scale numerical weather prediction with the cosmo
model: Description and sensitivities. *Monthly Weather Review* 139, 3887–3905.
- Belcher, S.E., Grant, A.L.M., Hanley, K.E., Fox-Kemper, B., Van Roekel, L., Sullivan,
P.P., Large, W.G., Brown, A., Hines, A., Calvert, D., Rutgersson, A., Pettersson,
H., Bidlot, J.R., Janssen, P.A.E.M., Polton, J.A., 2012. A global perspective on
725 Langmuir turbulence in the ocean surface boundary layer. *Geophys Res Lett* 39, 9.
doi:10.1029/2012GL052932.
- Breivik, Ø., Bidlot, J.R., Janssen, P.A., 2016. A stokes drift approximation based on
the phillips spectrum. *Ocean Modelling* 100, 49–56.
- Breivik, Ø., Janssen, P.A., Bidlot, J.R., 2014. Approximate stokes drift profiles in deep
730 water. *Journal of Physical Oceanography* 44, 2433–2445.
- Breivik, Ø., Mogensen, K., Bidlot, J.R., Balmaseda, M.A., Janssen, P.A., 2015. Surface
wave effects in the NEMO ocean model: Forced and coupled experiments. *Journal
of Geophysical Research: Oceans* 120, 2973–2992.
- Buckley, M.P., Veron, F., 2016. Structure of the airflow above surface waves. *Journal*
735 *of Physical Oceanography* 46, 1377–1397.
- Cai, Y., Wen, Y., Wu, L., Zhou, C., Zhang, F., 2017. Impact of wave breaking on
upper-ocean turbulence. *Journal of Geophysical Research: Oceans* 122, 1513–1528.
- Cavaleri, L., Abdalla, S., Benetazzo, A., Bertotti, L., Bidlot, J.R., Breivik, Ø., Carniel,
S., Jensen, R., Portilla-Yandun, J., Rogers, W., Roland, A., Sanchez-Arcilla, A.,
740 Smith, J., Staneva, J., Toledo, Y., van Vledder, G., van der Westhuysen, A., 2018.
Wave modelling in coastal and inner seas. *Prog Oceanogr* , 70doi:10.1016/j.
pcean.2018.03.010.
- Charnock, H., 1955. Wind stress on a water surface. *Quarterly Journal of the Royal
Meteorological Society* 81, 639–640.

- 745 Craig, P.D., Banner, M.L., 1994. Modeling wave-enhanced turbulence in the ocean surface layer. *Journal of Physical Oceanography* 24, 2546–2559.
- Davies, A.M., Kwong, S.C., Flather, R.A., 2000. On determining the role of wind wave turbulence and grid resolution upon computed storm driven currents. *Continental Shelf Research* 20, 1825–1888.
- 750 Davies, A.M., Lawrence, J., 1995. Modeling the effect of wave–current interaction on the three-dimensional wind-driven circulation of the eastern Irish sea. *Journal of Physical Oceanography* 25, 29–45.
- Dean, R.G., Dalrymple, R.A., 1991. *Water wave mechanics for engineers and scientists*. volume 2. world scientific publishing Co Inc.
- 755 Dietrich, J., Zijlema, M., Westerink, J., Holthuijsen, L., Dawson, C., Luettich, R.A., Jensen, R., Smith, J., Stelling, G., Stone, G., 2011. Modeling hurricane waves and storm surge using integrally-coupled, scalable computations. *Coastal Engineering* 58, 45–65.
- Drennan, W., Donelan, M., Terray, E., Katsaros, K., 1996. Oceanic turbulence dissipation measurements in swade. *Journal of Physical Oceanography* 26, 808–815.
- 760 Drennan, W.M., Kahma, K.K., Donelan, M.A., 1999. On momentum flux and velocity spectra over waves. *Boundary-Layer Meteorology* 92, 489–515.
- ECMWF, 2017. IFS Documentation CY43r1, Part VII: ECMWF Wave Model. ECMWF Model Documentation. European Centre for Medium-Range Weather Forecasts.
- 765 Edson, J.B., Jampana, V., Weller, R.A., Bigorre, S.P., Plueddemann, A.J., Fairall, C.W., Miller, S.D., Mahrt, L., Vickers, D., Hersbach, H., 2013. On the exchange of momentum over the open ocean. *Journal of Physical Oceanography* 43, 1589–1610.
- Egbert, G.D., Erofeeva, S.Y., 2002. Efficient Inverse Modeling of Barotropic Ocean Tides. *J Atmos Ocean Tech* 19, 183–204. doi:10.1175/1520-0426(2002)019<0183:EIMOBO>2.0.CO;2.

- Esters, L., Breivik, Ø., Landwehr, S., A ten Doeschate, G.S., Christensen, K., Bidlot, J.R., Ward, B., 2018. Turbulence Scaling Comparisons in the Ocean Surface Boundary Layer. *J Geophys Res: Oceans*, 20doi:10.1002/2017JC013525.
- 775 Fan, Y., Griffies, S.M., 2014. Impacts of parameterized langmuir turbulence and non-breaking wave mixing in global climate simulations. *Journal of Climate* 27, 4752–4775.
- Fan, Y., Yu, Z., 2017. Surface gravity wave effect in turbulent kinetic energy flux across the air-sea interface, in: *EGU General Assembly Conference Abstracts*, p. 5624.
- 780 Feddersen, F., Trowbridge, J.H., Williams III, A., 2007. Vertical structure of dissipation in the nearshore. *Journal of Physical Oceanography* 37, 1764–1777.
- Gerbi, G.P., Trowbridge, J.H., Terray, E.A., Plueddemann, A.J., Kukulka, T., 2009. Observations of turbulence in the ocean surface boundary layer: Energetics and transport. *Journal of Physical Oceanography* 39, 1077–1096.
- 785 Guan, C., Xie, L., 2004. On the linear parameterization of drag coefficient over sea surface. *Journal of Physical Oceanography* 34, 2847–2851.
- Harcourt, R., 2013. A Second-Moment Closure Model of Langmuir Turbulence. *J Phys Oceanogr* 43, 673–697. doi:10.1175/JPO-D-12-0105.1.
- Harcourt, R.R., 2015. An improved second-moment closure model of langmuir turbulence. *Journal of Physical Oceanography* 45, 84–103.
- 790 Harcourt, R.R., D’Asaro, E.A., 2008. Large-eddy simulation of Langmuir turbulence in pure wind seas. *J Phys Oceanogr* 38, 1542–1562.
- Hasselmann, K., 1970. Wave-driven inertial oscillations. *Geophysical and Astrophysical Fluid Dynamics* 1, 463–502.
- 795 Hasselmann, K., 1991. Ocean circulation and climate change. *Tellus A: Dynamic Meteorology and Oceanography* 43, 82–103.

- Hasselmann, S., Hasselmann, K., Allender, J., Barnett, T., 1985. Computations and parameterizations of the nonlinear energy transfer in a gravity-wave spectrum. part ii: Parameterizations of the nonlinear energy transfer for application in wave models. *Journal of Physical Oceanography* 15, 1378–1391.
- He, H., Chen, D., 2011. Effects of surface wave breaking on the oceanic boundary layer. *Geophysical Research Letters* 38.
- Hersbach, H., 2011. Sea Surface Roughness and Drag Coefficient as Functions of Neutral Wind Speed 41, 247–251. doi:10.1175/2010JPO4567.1.
- Högström, U., Rutgersson, A., Sahlée, E., Smedman, A.S., Hristov, T.S., Drennan, W., Kahma, K., 2013. Air–sea interaction features in the baltic sea and at a pacific trade-wind site: an inter-comparison study. *Boundary-layer meteorology* 147, 139–163.
- Högström, U., Sahlée, E., Smedman, A.S., Rutgersson, A., Nilsson, E., Kahma, K.K., Drennan, W.M., 2015. Surface stress over the ocean in swell-dominated conditions during moderate winds. *Journal of the Atmospheric Sciences* 72, 4777–4795.
- Huang, C.J., Qiao, F., 2010. Wave-turbulence interaction and its induced mixing in the upper ocean. *Journal of Geophysical Research: Oceans* 115.
- Janssen, F., Schrum, C., Backhaus, J.O., 1999. A climatological data set of temperature and salinity for the baltic sea and the north sea. *Deutsche Hydrografische Zeitschrift* 51, 5.
- Janssen, P., 2004. The interaction of ocean waves and wind. Cambridge University Press, Cambridge, UK.
- Janssen, P., 2012. Ocean Wave Effects on the Daily Cycle in SST. *J Geophys Res: Oceans* 117, 24. doi:10/mth.
- Janssen, P.A., 1989. Wave-induced stress and the drag of air flow over sea waves. *Journal of Physical Oceanography* 19, 745–754.
- Janssen, P.A., 1991. Quasi-linear theory of wind-wave generation applied to wave forecasting. *Journal of Physical Oceanography* 21, 1631–1642.

- 825 Jones, J., Davies, A., 1998. Storm surge computations for the Irish sea using a three-dimensional numerical model including wave–current interaction. *Continental Shelf Research* 18, 201–251.
- Jones, N.L., Monismith, S.G., 2008. The influence of whitecapping waves on the vertical structure of turbulence in a shallow estuarine embayment. *Journal of Physical Oceanography* 38, 1563–1580.
- 830 Kantha, L.H., Clayson, C.A., 2004. On the effect of surface gravity waves on mixing in the oceanic mixed layer. *Ocean Modelling* 6, 101–124.
- Kenyon, K.E., 1969. Stokes drift for random gravity waves. *Journal of Geophysical Research* 74, 6991–6994.
- Komen, G.J., Cavaleri, L., Donelan, M., Hasselmann, K., Hasselmann, S., Janssen, P.A.E.M., 1994. *Dynamics and Modelling of Ocean Waves*. Cambridge University Press, Cambridge.
- 835 Large, W.G., McWilliams, J.C., Doney, S.C., 1994. Oceanic vertical mixing: A review and a model with a nonlocal boundary layer parameterization. *Rev Geophys* 32, 363–403. doi:10.1029/94RG01872.
- 840 Large, W.G., Yeager, S., 2009. The global climatology of an interannually varying air–sea flux data set. *Climate dynamics* 33, 341–364.
- Large, W.G., Yeager, S.G., 2004. Diurnal to decadal global forcing for ocean and sea-ice models: the data sets and flux climatologies. *National Center for Atmospheric Research Boulder*.
- 845 Lehmann, A., Myrberg, K., 2008. Upwelling in the Baltic sea: a review. *Journal of Marine Systems* 74, S3–S12.
- Lehmann, A., Myrberg, K., Höflich, K., 2012. A statistical approach to coastal upwelling in the Baltic sea based on the analysis of satellite data for 1990–2009. *Oceanologia* 54, 369–393.

- 850 Li, M., Garrett, C., 1993. Cell merging and the jet/downwelling ratio in langmuir circulation. *Journal of Marine Research* 51, 737–769.
- Li, Q., Fox-Kemper, B., Breivik, Ø., Webb, A., 2017. Statistical Models of Global Langmuir Mixing. *Ocean Model* 113, 95–114. doi:10.1016/j.ocemod.2017.03.016.
- 855 Li, Q., Webb, A., Fox-Kemper, B., Craig, A., Danabasoglu, G., Large, W.G., Vertenstein, M., 2016. Langmuir mixing effects on global climate: WAVEWATCH III in CESM. *Ocean Model* 103, 145–160. doi:10.1016/j.ocemod.2015.07.020.
- Lips, I., Lips, U., 2008. Abiotic factors influencing cyanobacterial bloom development in the gulf of finland (baltic sea). *Hydrobiologia* 614, 133–140.
- 860 Madec, G., et al., 2015. Nemo ocean engine .
- Mastenbroek, C., Burgers, G., Janssen, P., 1993. The dynamical coupling of a wave model and a storm surge model through the atmospheric boundary layer. *J Phys Oceanogr* 23, 1856–1866. doi:10/ccb4f9.
- McWilliams, J.C., Fox-Kemper, B., 2013. Oceanic wave-balanced surface fronts and
865 filaments. *Journal of Fluid Mechanics* 730, 464–490.
- McWilliams, J.C., Gula, J., Molemaker, M.J., Renault, L., Shchepetkin, A.F., 2015. Filament frontogenesis by boundary layer turbulence. *Journal of Physical Oceanography* 45, 1988–2005.
- McWilliams, J.C., Huckle, E., Liang, J., Sullivan, P.P., 2014. Langmuir turbulence in
870 swell. *Journal of Physical Oceanography* 44, 870–890.
- McWilliams, J.C., Huckle, E., Liang, J.H., Sullivan, P.P., 2012. The wavy ekman layer: Langmuir circulations, breaking waves, and reynolds stress. *Journal of Physical Oceanography* 42, 1793–1816.
- McWilliams, J.C., Restrepo, J.M., 1999. The Wave-driven Ocean Circulation. *J Phys Oceanogr* 29, 2523–2540. doi:10/dwj9tj.
- 875

- McWilliams, J.C., Sullivan, P.P., 2000. Vertical mixing by langmuir circulations. *Spill Science & Technology Bulletin* 6, 225–237.
- Mellor, G.L., Yamada, T., 1982. Development of a turbulent closure model for geophysical fluid problems. *Rev Geophys Space Phys* 20, 851–875. doi:10.1029/RG020i004p00851.
- 880 Mogensen, K.S., Magnusson, L., Bidlot, J.R., 2017. Tropical cyclone sensitivity to ocean coupling in the ECMWF coupled model. *J Geophys Res: Oceans* 122, 4392–4412. doi:10.1002/2017JC012753.
- Norman, M., Parampil, S.R., Rutgersson, A., Sahlée, E., 2013. Influence of coastal upwelling on the air–sea gas exchange of co₂ in a baltic sea basin. *Tellus B: Chemical and Physical Meteorology* 65, 21831.
- 885 Phillips, O.M., 1958. The equilibrium range in the spectrum of wind-generated waves. *J Fluid Mech* 4, 426–434. doi:10.1017/S0022112058000550.
- Pierson, Jr, W.J., Moskowitz, L., 1964. A proposed spectral form for fully developed wind seas based on the similarity theory of S A Kitaigorodskii. *J Geophys Res* 69, 5181–5190.
- 890 Polton, J.A., 2009. A wave averaged energy equation: Comment on global estimates of wind energy input to subinertial motions in the ekman-stokes layer by bin liu, kejian wu and changlong guan. *Journal of oceanography* 65, 665–668.
- Polton, J.A., Lewis, D.M., Belcher, S.E., 2005. The role of wave-induced coriolis–stokes forcing on the wind-driven mixed layer. *Journal of Physical Oceanography* 35, 444–457.
- 895 Powell, M.D., Vickery, P.J., Reinhold, T.A., 2003. Reduced drag coefficient for high wind speeds in tropical cyclones. *Nature* 422, 279–283.
- Qiao, F., Yuan, Y., Yang, Y., Zheng, Q., Xia, C., Ma, J., 2004. Wave-induced mixing in the upper ocean: Distribution and application to a global ocean circulation model. *Geophysical Research Letters* 31.
- 900

- Reichl, B.G., Wang, D., Hara, T., Ginis, I., Kukulka, T., 2016. Langmuir turbulence parameterization in tropical cyclone conditions. *Journal of Physical Oceanography* 46, 863–886.
- 905
- Rutgersson, A., Smedman, A.S., Högström, U., 2001. Use of conventional stability parameters during swell. *Journal of Geophysical Research: Oceans* 106, 27117–27134.
- Saetra, Ø., Albrechtsen, J., Janssen, P.A., 2007. Sea-state-dependent momentum fluxes for ocean modeling. *Journal of Physical Oceanography* 37, 2714–2725.
- 910
- Sheldon, L., Czaja, A., 2014. Seasonal and interannual variability of an index of deep atmospheric convection over western boundary currents. *Q J R Meteorol Soc* 140, 22–30. doi:10.1002/qj.2103.
- Smyth, W.D., Skillingstad, E.D., Crawford, G.B., Wijesekera, H., 2002. Nonlocal fluxes and stokes drift effects in the k-profile parameterization. *Ocean Dynamics* 52, 104–115.
- 915
- Sproson, D., Sahlée, E., 2014. Modelling the impact of baltic sea upwelling on the atmospheric boundary layer. *Tellus A: Dynamic Meteorology and Oceanography* 66, 24041.
- Staneva, J., Alari, V., Breivik, Ø., Bidlot, J.R., Mogensen, K., 2017. Effects of wave-induced forcing on a circulation model of the north sea. *Ocean Dynamics* 67, 81–101.
- 920
- Stokes, G.G., 1847. On the theory of oscillatory waves. *Trans Cambridge Philos Soc* 8, 441–473.
- Sullivan, P.P., McWilliams, J.C., 2010. Dynamics of winds and currents coupled to surface waves. *Annual Review of Fluid Mechanics* 42.
- 925
- Sullivan, P.P., McWILLIAMS, J.C., Melville, W.K., 2004. The oceanic boundary layer driven by wave breaking with stochastic variability. part 1. direct numerical simulations. *Journal of Fluid Mechanics* 507, 143–174.

- 930 Sullivan, P.P., McWILLIAMS, J.C., Melville, W.K., 2007. Surface gravity wave effects in the oceanic boundary layer: Large-eddy simulation with vortex force and stochastic breakers. *Journal of Fluid Mechanics* 593, 405–452.
- Sutherland, P., Melville, W.K., 2015. Field measurements of surface and near-surface turbulence in the presence of breaking waves. *Journal of Physical Oceanography* 45, 943–965.
935
- Suzuki, N., Fox-Kemper, B., 2016. Understanding stokes forces in the wave-averaged equations. *Journal of Geophysical Research: Oceans* 121, 3579–3596.
- Takaya, Y., Bidlot, J.R., Beljaars, A., Janssen, P., 2010. Refinements to a prognostic scheme of skin sea surface temperature. *J Geophys Res: Oceans* 115. doi:10.1029/2009JC005985.
940
- Terray, E., Donelan, M., Agrawal, Y., Drennan, W., Kahma, K., Williams, A., Hwang, P., Kitaigorodskii, S., 1996. Estimates of kinetic energy dissipation under breaking waves. *Journal of Physical Oceanography* 26, 792–807.
- Thomson, J., Schwendeman, M.S., Zippel, S.F., Moghimi, S., Gemmrich, J., Rogers, W.E., 2016. Wave-breaking turbulence in the ocean surface layer. *Journal of Physical Oceanography* 46, 1857–1870.
945
- Umlauf, L., Burchard, H., 2003. A generic length-scale equation for geophysical turbulence models. *J Mar Res* 61, 235–265. doi:10.1357/002224003322005087.
- Umlauf, L., Burchard, H., 2005. Second-order turbulence closure models for geophysical boundary layers. A review of recent work. *Cont Shelf Res* 25, 795–827.
950 doi:10.1016/j.csr.2004.08.004.
- Van Roekel, L.P., Fox-Kemper, B., Sullivan, P.P., Hamlington, P.E., Haney, S.R., 2012. The form and orientation of Langmuir cells for misaligned winds and waves. *J Geophys Res: Oceans* 117, 22, C05001. doi:10.1029/2011JC007516.
- 955 Veron, F., Melville, W.K., Lenain, L., 2008. Wave-coherent air–sea heat flux. *Journal of Physical Oceanography* 38, 788–802.

- Wahle, K., Staneva, J., Koch, W., Fenoglio-Marc, L., Ho-Hagemann, H.T., Stanev, E.V., 2017. An atmosphere–wave regional coupled model: improving predictions of wave heights and surface winds in the southern north sea. *Ocean Science* 13, 289.
- 960 WAMDI, 1988. The wam modela third generation ocean wave prediction model. *Journal of Physical Oceanography* 18, 1775–1810.
- Wang, J., Shen, Y., 2011. Development and validation of a three-dimensional, wave-current coupled model on unstructured meshes. *Science China Physics, Mechanics and Astronomy* 54, 42–58.
- 965 Wu, L., Rutgersson, A., Nilsson, E., 2017. Atmospheric boundary layer turbulence closure scheme for wind-following swell conditions. *Journal of the Atmospheric Sciences* 74, 2363–2382.
- Wu, L., Rutgersson, A., Sahlée, E., 2015. Upper-ocean mixing due to surface gravity waves. *Journal of Geophysical Research: Oceans* 120, 8210–8228.
- 970 Wu, L., Rutgersson, A., Sahlée, E., Guo Larsén, X., 2016. Swell impact on wind stress and atmospheric mixing in a regional coupled atmosphere-wave model. *Journal of Geophysical Research: Oceans* 121, 4633–4648.

Table 1: Experiment design.

Experiments	Group
CTL	Control simulation, without wave model
STCOR	CTL + Stokes-Coriolis forcing
STFUL	STCOR + Stokes drift impact on tracer and mass advection
FULL	STFUL + sea-state-dependent momentum and energy fluxes

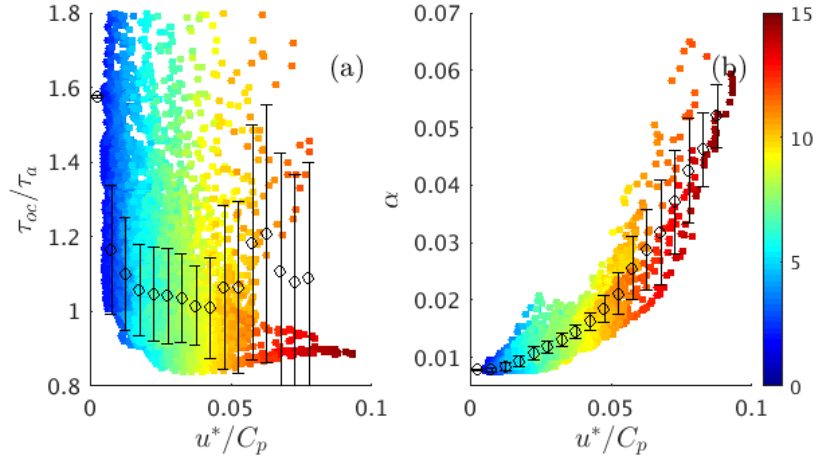


Figure 1: The relationship between u^*/C_p and the normalised wind stress τ_{oc}/τ_a (a), and the Charnock coefficient α (b). The color scale represents U_{10} [m s^{-1}]. The data are from the hourly output data from wave model simulation during December 7th, 2015 in the blue box domain shown in Fig. 2. The circles show the bin average with standard deviation shown as bars. The vertical axis is limited for legibility as only a few data points exceed 1.8.

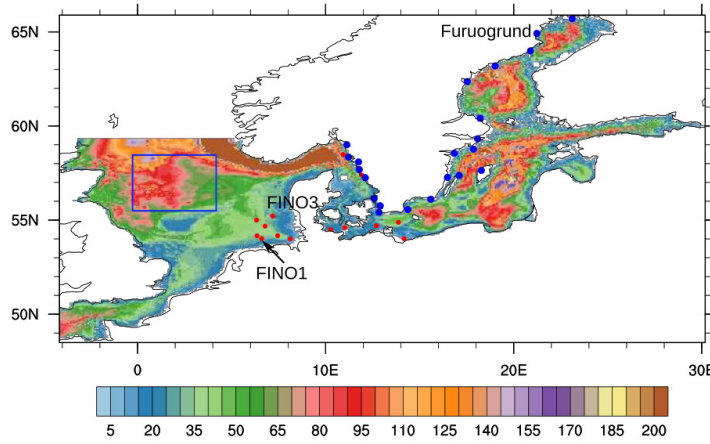


Figure 2: Topography of the model domain and the buoy locations used in this study [m]. The red circles represent temperature measurement stations and the blue dots represent sea level stations used in this study. The blue box represents the area in which data are used in Fig. 18 and 1.

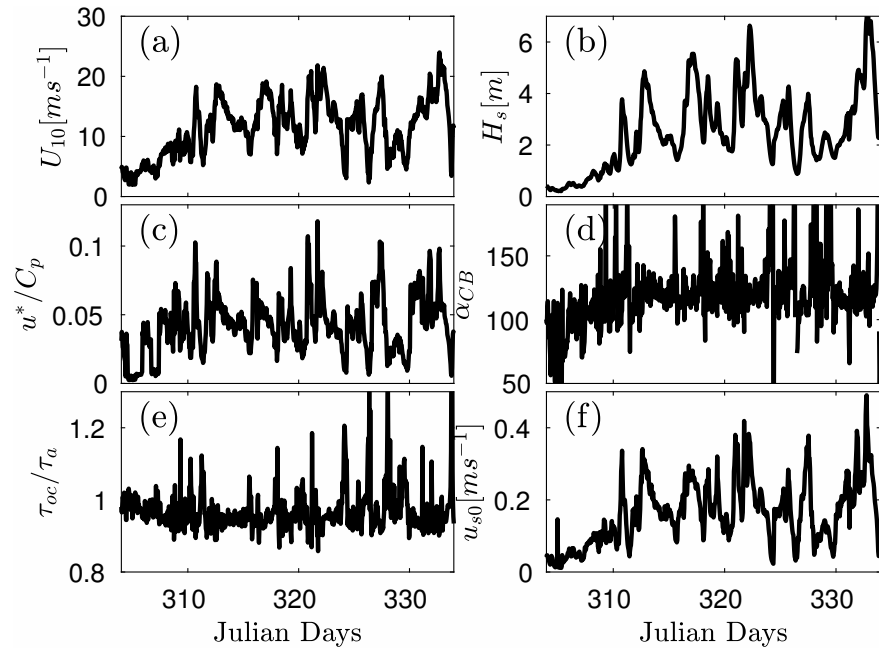


Figure 3: The time series of the mean wind speed U_{10} (a) , (b) the significant wave height, (c) inverse wave age u^*/C_p , (d) α_{CB} , (e) τ_{oc}/τ_a and the surface Stokes drift u_{s0} at the FINO1 station during the storm simulation period. The wind and wave data are from DWD.

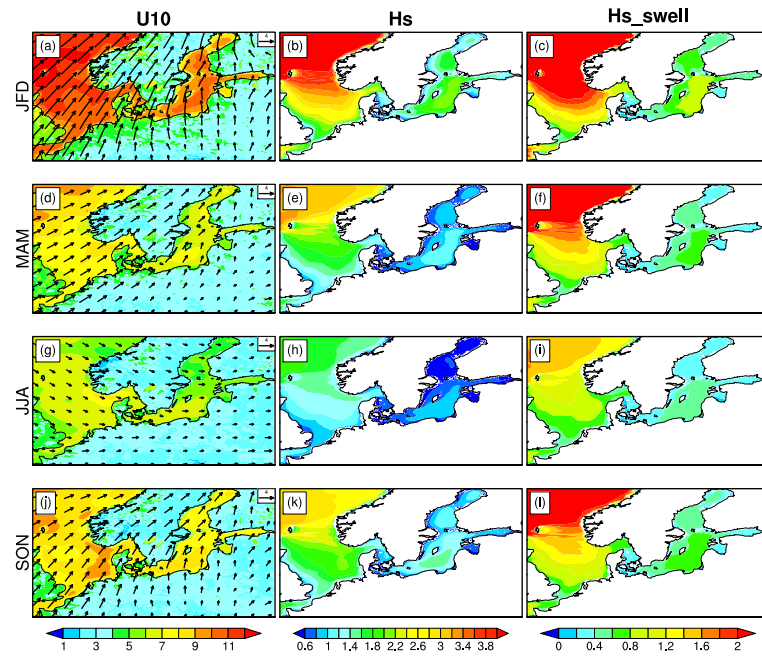


Figure 4: The mean forcing field (from the left to right columns) for wind speed at 10 m [m s^{-1}], significant wave height [m], and the swell significant wave height [m]. The four rows show the mean value in JFD, MAM, JJA, and SON.

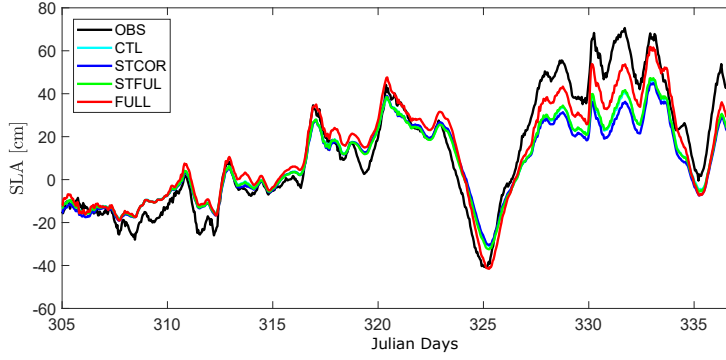


Figure 5: The simulation results of SLA from the four experiments and the measurements during the storm period at station Furuogrund. The lines of STFUL almost overlap CTL.

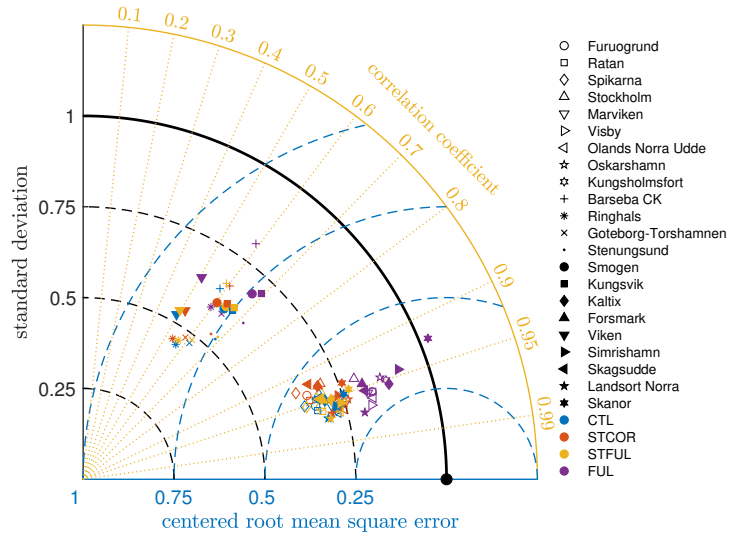


Figure 6: The Taylor diagram of SLA for all stations.

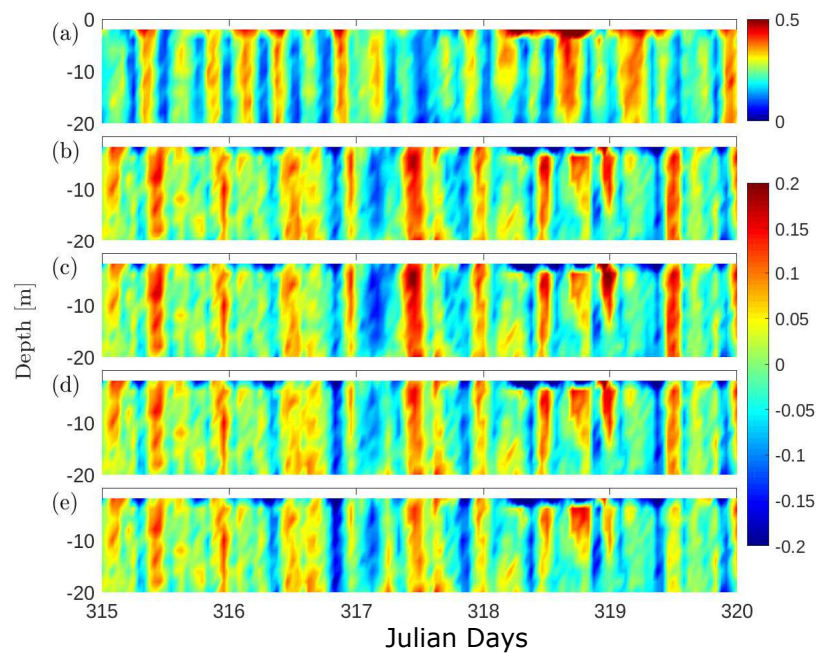


Figure 7: The velocity profiles at FINO3 station for Julian days 315-320 in year 2015: (a) hourly mean measurements (OBS), hourly mean difference CTL-OBS (b), STCOR-OBS (c), STFUL-OBS (d), and FULL-OBS (e).

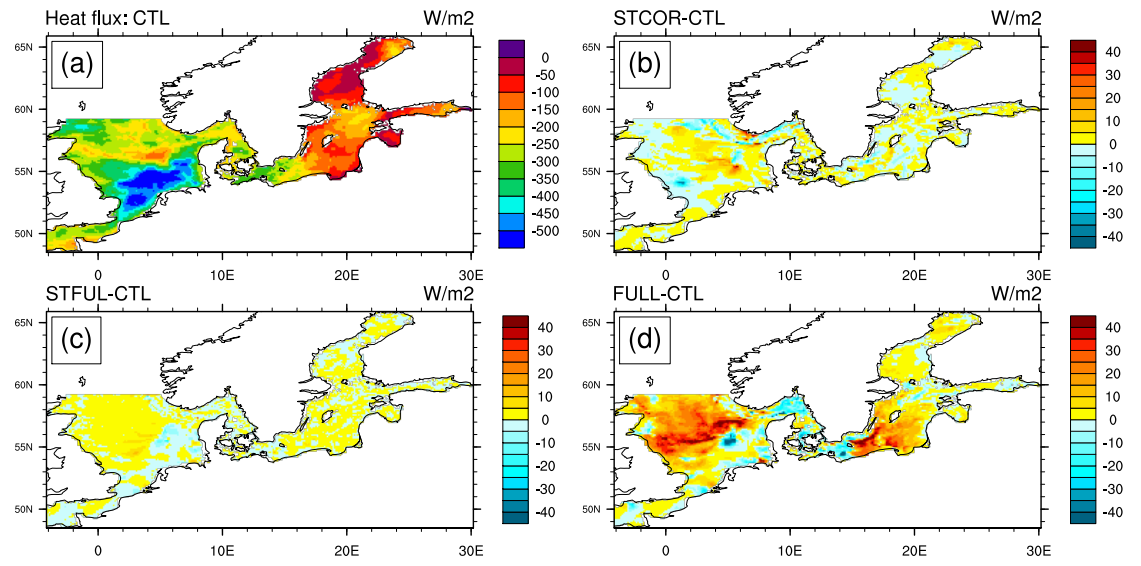


Figure 8: The net heat flux [W m^{-2}] at 01:00 UTC, Julian day 318 (November 14) for CTL (a), the difference STCOR-CTL (b), STFUL-CTL (c), and FULL-CTL (d). The negative flux in (a) means heat flux from the ocean to the atmosphere. The positive flux in (b)-(d) means that the wave influence reduces the flux from the ocean to the atmosphere.

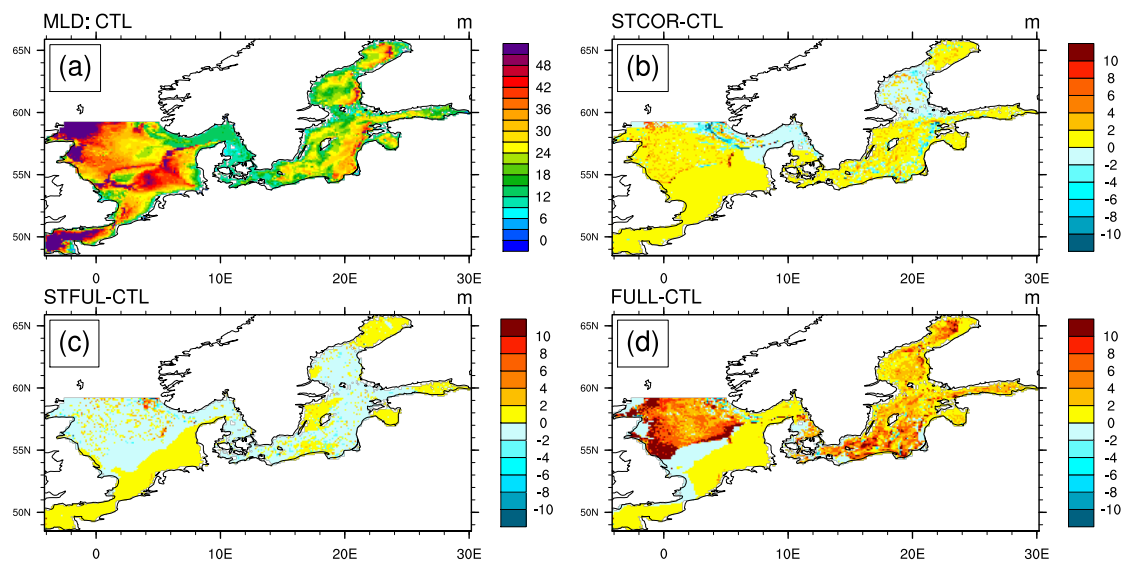


Figure 9: The mixed layer depth [m] at 01:00 UTC, Julian day 318 (November 14) for CTL (a), the difference STCOR-CTL (b), STFUL-CTL (c), and FULL-CTL (d).

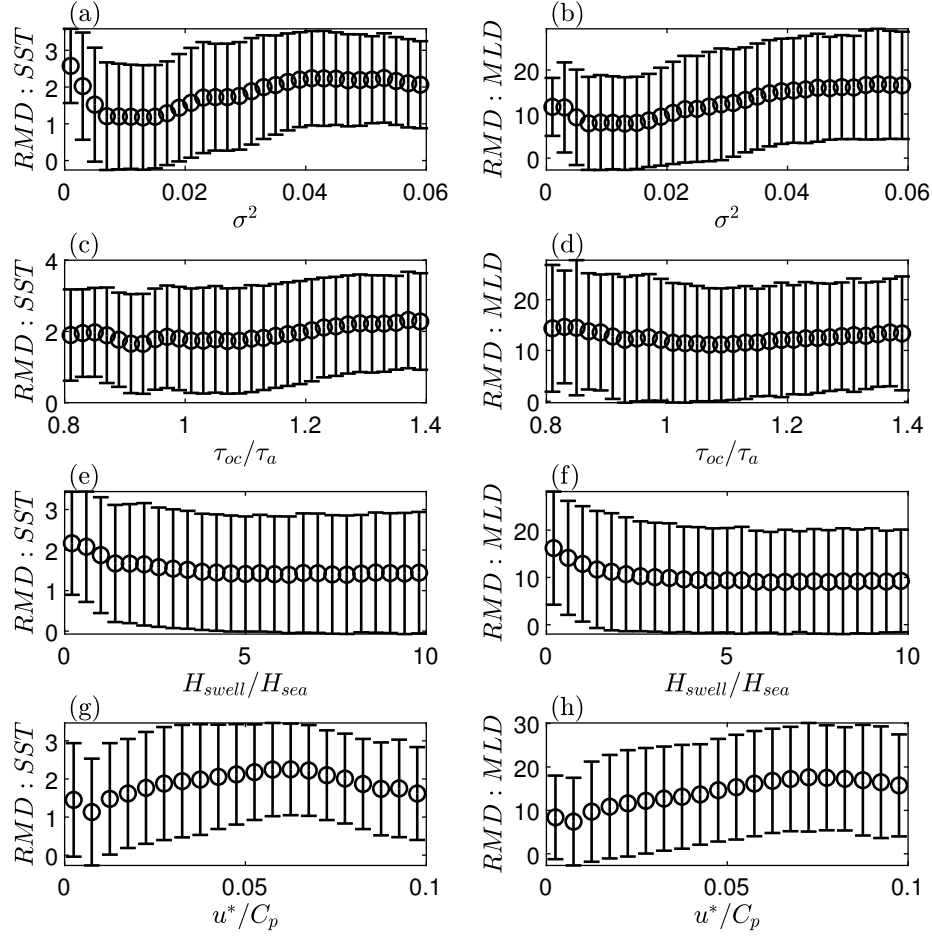


Figure 10: The relative mean difference of SST [%] under different mean square wave slope (a), (c) τ_{oc}/τ_a , (e) H_{swell}/H_{sea} , and (g) inverse wave age u^*/C_p . The relative mean difference of the mixed layer depth under different mean square wave slope (b), (d) τ_{oc}/τ_a , (f) H_{swell}/H_{sea} , and (h) inverse wave age u^*/C_p . The data used in this figure are the hourly data from the area shown in blue box in Fig. 2.

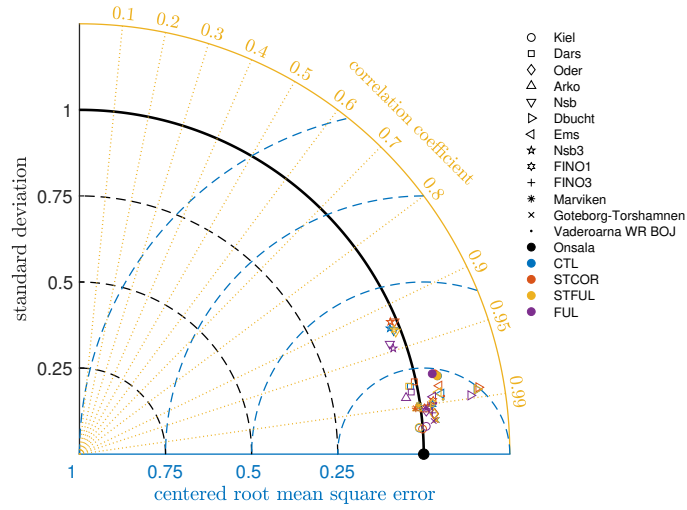


Figure 11: The Taylor diagram of SST for all stations.

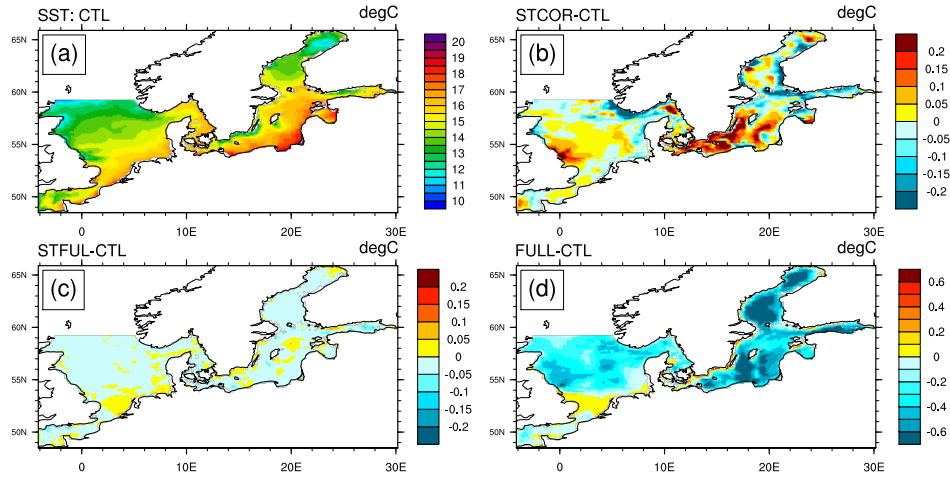


Figure 12: The mean SST [$^{\circ}\text{C}$] of June, July and August in 2015 (a), the mean SST difference STCOR-CTL (b), STFUL-CTL (c), and FULL-CTL (d) for the months June, July and August.

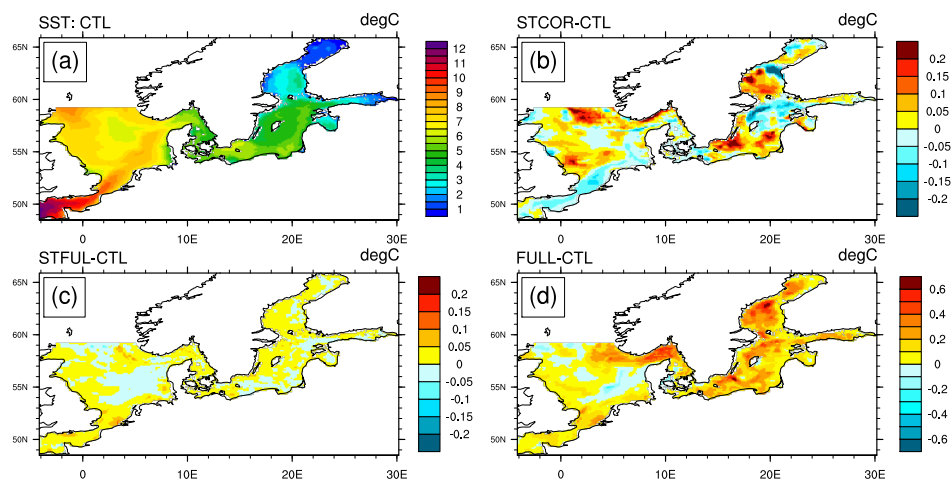


Figure 13: Same as Fig. 12, but for the months December, January and February 2015.

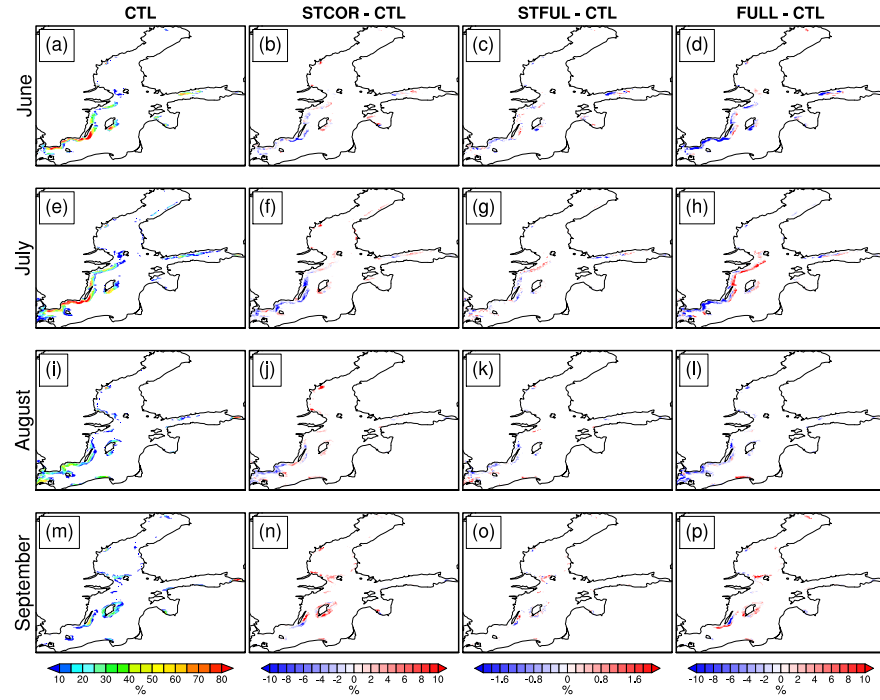


Figure 14: The frequency of upwelling in the control experiment for the months of June, July, August and September, 2015 [%]. The four columns represent the upwelling frequency in the control experiment, the difference between the experiment STCOR, STFUL, FULL and the control experiments. The four rows represent the months June, July, August and September 2015. The criterion for upwelling are fulfilled when the SST difference from the zonal mean temperature is greater than 2.5 °C.

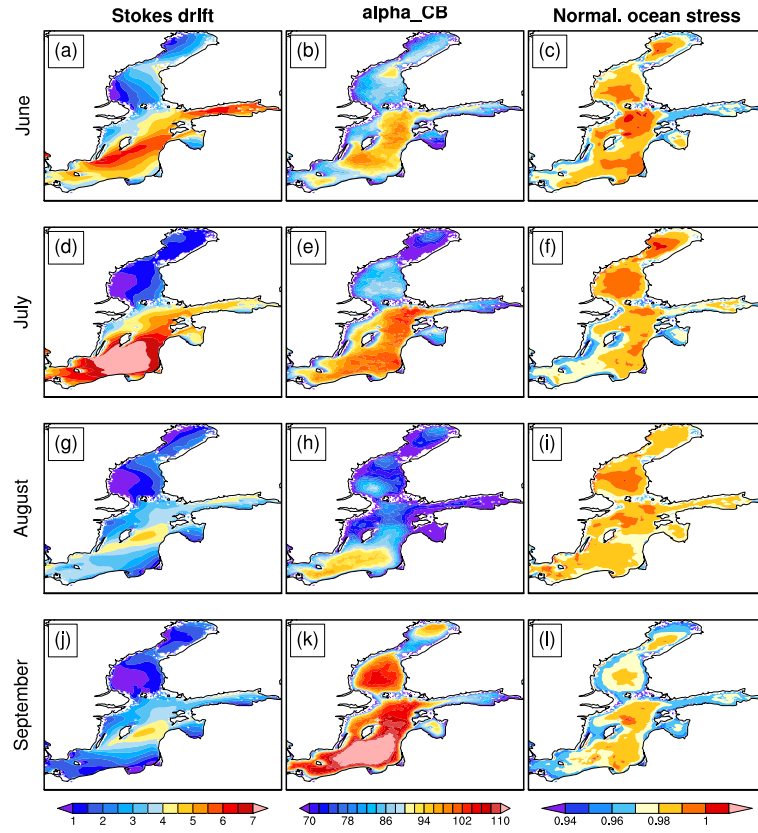


Figure 15: The three columns show the mean surface Stokes drift, the wave energy coefficient α_{CB} , and the normalized ocean stress τ_{oc}/τ_a . The four rows are data for June, July, August, and September 2015.

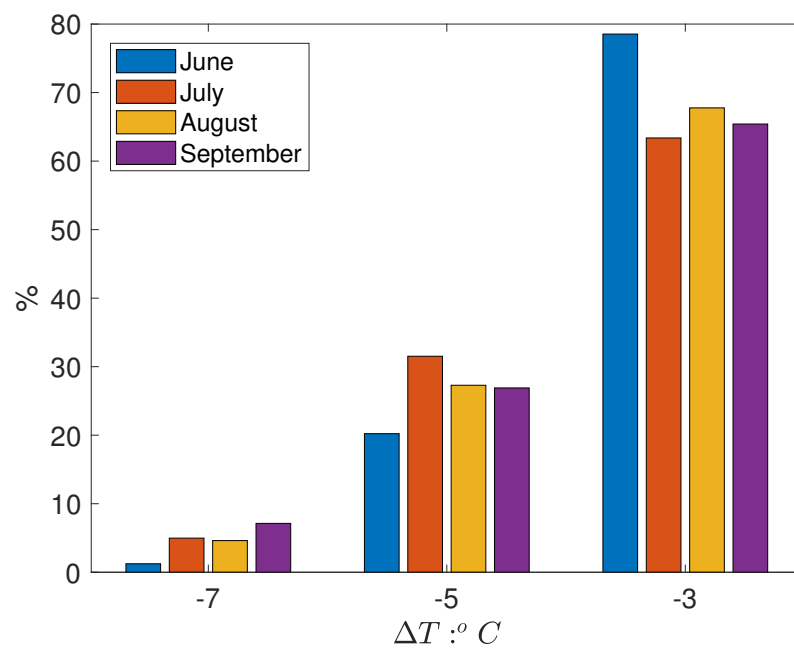


Figure 16: The normalized distribution of the upwelling intensity (with 2 $^{\circ}\text{C}$ bins) in June, July, August and September 2015.

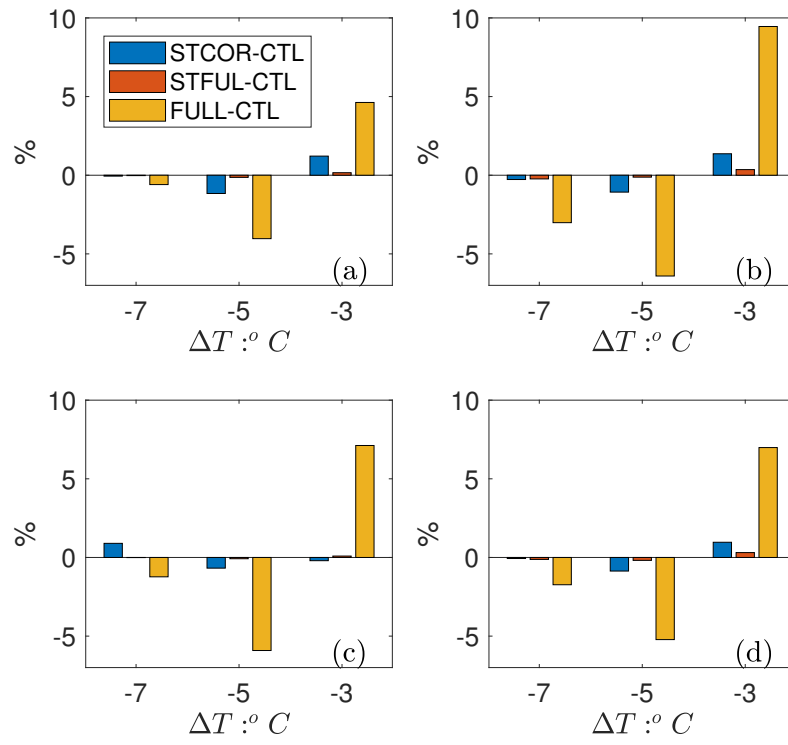


Figure 17: The wave effect on the distribution of the upwelling intensity in (a) June, (b) July, (c) August, and (d) September 2015.

Table 2: The location of the stations shown in Fig. 2.

SLA Station	Lat	Lon	SST Station	Lat	Lon
Furuogrund	64.92	21.23	Kiel	54.50	10.27
Ratan	63.99	20.89	Dars	54.70	12.70
Spikarna	62.36	17.53	Oder	54.01	14.17
Stockholm	59.32	18.08	Arko	54.88	13.87
Marviken	58.55	16.84	Nsb	55.00	6.33
Visby	57.64	18.28	Dbucht	55.17	7.45
Olands Norra Udde	57.37	17.10	Ems	54.17	6.35
Oskarshamn	57.27	16.48	Nsb3	54.68	6.78
Kungsholmsfort	56.11	15.59	FINO1	54.02	6.60
Barseba CK	55.76	12.90	FINO3	55.22	7.18
Ringhals	57.25	12.11	Marviken	58.55	16.84
Goteborg-Torshamnen	57.68	11.79	Goteborg-Torshamnen	57.68	11.79
Stenungsund	58.09	11.83	Vaderoarna WR BOJ	58.48	10.93
Smogen	58.35	11.22	Onsala	57.39	11.92
Kungsvik	59.00	11.13			
Kaltix	65.70	23.10			
Forsmark	60.41	18.21			
Viken	56.14	12.58			
Simrishamn	55.56	14.36			
Skagsudde	63.19	19.01			
Landsort Norra	58.77	17.86			
Skanor	55.42	12.83			

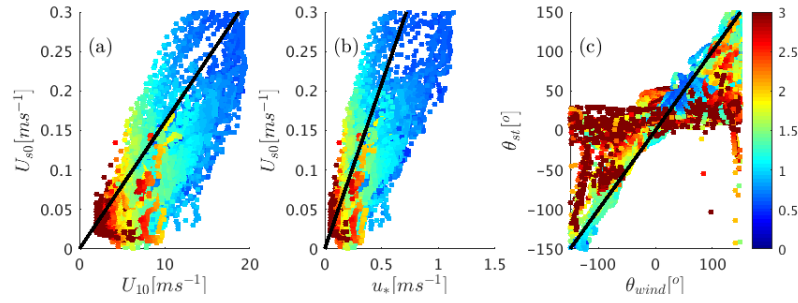


Figure 18: The relationship between U_{10} and the magnitude of the surface Stokes drift is shown in (a) and the black line represents the $u_{s0} = 0.016U_{10}$ (Li and Garrett, 1993). Panel (b) shows the relationship between the u_* and the magnitude of the surface Stokes drift, and the black line shows $u_{s0} = 0.377|\tau/\rho_w|^{1/2}$ (Maded et al., 2015). Panel (c) shows the surface Stokes drift direction and the direction of U_{10} . Here the black line represents the 1:1 line. The color represents the wave age, C_p/U_{10} . The data are from the hourly output data from wave model simulation during Julian day 341 (December 7th), 2015 in the blue box domain shown in Fig. 2.